**14. Black Hole Entropy as TD Entropy** 

- Prunkl & Timpson (2019)

# 1. Motivations

 $S_{\rm BH} = A/4$ 

- The surface area A of the event horizon of the black hole.

 $S_{\rm TD}(\sigma_2) \equiv \int_{\sigma_1}^{\sigma_2} \frac{\delta Q_R}{T} + S_0 \quad \longleftarrow$ 

The ratio of the change in heat to temperature of a reversible process that connects an initial state to a final state.

Is  $S_{\rm BH} = S_{\rm TD}$ ?

Is the thermodynamic entropy of a black hole proportional to the surface area of its event horizon?

## Bekenstein's motivations:

- The laws of black hole mechanics are formally *Only establishes an analogy.* similar to the laws of thermodynamics.
- If we don't identify the TD entropy of a black hole with its event horizon, then the 2<sup>nd</sup> Law is violated.

### What's needed:

An explicit demonstration that a black hole can undergo a reversible process.

- 1. Motivations
- 2. Problem: Negative Heat Capacity
- 3. A Black Hole Carnot Cycle

Only persuasive for

reversible processes.

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## 2. Problem: Negative Heat Capacity

- <u>*Recall*</u>: A reversible process is a process that takes an initial equilibrium state to a final equilibrium state.
- <u>So</u>: If a black hole can undergo a reversible process (and thus be ascribed  $S_{TD}$ ), it must be able to be in an equilibrium state with its environment.

## <u>Problem</u>: A black hole has a negative heat capacity!

Heat capacity C = amount of absorbed heat  $\delta Q$  needed to change temperature by dT:  $C = \delta Q/dT$ 

<u>*Claim*</u>: If  $S_{BH} = S_{TD}$  and  $T = (1/2\pi)\kappa$ , then a black hole has negative *C*.

<u>*Proof*</u>: For a Schwarzschild black hole, R = 2M,  $A = 4\pi R^2$ ,  $\kappa = 1/4M$ 

- <u>So</u>:  $S_{\rm BH} = A/4 = 4\pi M^2$  or  $dS_{\rm BH} = 8\pi M dM$
- <u>And</u>:  $T = (1/2\pi)\kappa = 1/8\pi M$  or  $dT = -(1/8\pi M^2)dM$
- <u>Suppose</u>:  $S_{BH} = S_{TD}$
- <u>Then</u>:  $\delta Q = T dS_{\text{TD}} = T dS_{\text{BH}} = dM$
- <u>Thus</u>:  $C = \delta Q/dT = -8\pi M^2$

## 2. Problem: Negative Heat Capacity

- <u>*Recall*</u>: A reversible process is a process that takes an initial equilibrium state to a final equilibrium state.
- <u>So</u>: If a black hole can undergo a reversible process (and thus be ascribed  $S_{TD}$ ), it must be able to be in an equilibrium state with its environment.

# <u>Problem</u>: A black hole has a negative heat capacity!

Heat capacity C = amount of absorbed heat  $\delta Q$  needed to change temperature by dT:  $C = \delta Q/dT$ 

- If *C* is negative, then when heat is absorbed, temperature *decreases!*
- <u>*This is a problem*</u>: How can a black hole be in a stable equilibrium state with another physical system with positive heat capacity?

/V As that system cools, it emits heat, which is absorbed by the black hole, which also cools as a result.

The black hole and the other system can never reach an equilibrium state! -  $S_{\rm BH}$  must be statistical mechanical microcanonical entropy.

"The fact that the temperature of a black hole decreases as the mass increases means that black holes cannot be in stable thermal equilibrium in the situations in which there is an indefinitely large amount of energy available... [T]his implies that the normal statistical mechanical canonical ensemble cannot be applied to gravitating systems. Instead one has to use microcanonical ensembles in which one considers all the possible configurations of a system with a given energy."



"A black hole of given mass, angular momentum, and charge can have a large number of different unobservable internal configurations which reflect the possible internal configurations of the matter which collapsed to produce the hole. The logarithm of this number can be regarded as the entropy of the black hole and is a measure of the amount of information about the initial state which was lost in the formation of the black hole."



$$S_{\rm BH} = S_{\rm Gibbs}(\rho_{\rm mc}) = k \ln \Omega(E)$$

The number of microstates of the black hole?
 The information "about the initial state which was lost in the formation of the black hole"?

## Prunkl & Timpson's response:

 "Even though black holes cannot be in stable equilibrium with an infinite heat bath, they can be in stable equilibrium with a photon gas and enclosed in a box, for a certain range of parameters."

<u>Specific claim</u>: For a black hole of mass *M* in a photon gas of volume *V*, there are values of *M* and *V* that make the equilibrium condition  $T_{BH} = T_{gas}$  stable against fluctuations in temperature, in the sense:



"The gas can react quickly enough to any fluctuations in the black hole."



Black hole of mass M in photon gas of volume V





Black hole of mass M in photon gas of volume V

Energy of photon gas from Planck's Law.

• <u>*First*</u>: Rewrite (1) in terms of *M*:  $E_{\text{total}} = E_{\text{BH}} + E_{\text{gas}} = M + \frac{\pi^2}{15} V T_{\text{gas}}^4$ 

- Now solve for  $T_{gas}$ :

$$T_{\text{gas}} = \left[\frac{E_{\text{total}} - M}{\frac{\pi^2}{15}V}\right]^{\frac{1}{4}}$$

 $\left[\frac{E_{\text{total}} - M}{\frac{\pi^2}{V}}\right]^{74} = \frac{1}{8\pi M} \quad \Leftarrow \text{Assumes a black hole has a temperature given by}$ - Now substitute into (1) with  $T_{\rm BH} = 1/8\pi M$ :

Hawking radiation.

- Or: 
$$M^4(M - E_{\text{total}}) + \beta V = 0$$
,  $\beta = \frac{1}{15(8)^4 \pi^2}$ 

To solve this for M, construct the graph of the function  $f(M) = M^4(M - E_{\text{total}}) + \beta V$ and see where it intersects the M axis...

- Let  $f(M) = M^4(M E_{\text{total}}) + \beta V$   $f(M) = 0 \Leftrightarrow T_{\text{BH}} = T_{\text{gas}}$ 
  - Max/min occur at values of *M* for which df/dM = 0:

 $df/dM = 5M^4 - 4M^3E_{\text{total}} = 0$  for M = 0 and  $M = \frac{4}{5}E_{\text{total}}$ 

- <u>Note</u>:  $f(M) \to -\infty$  as  $M \to -\infty$  and  $f(M) \to +\infty$  as  $M \to +\infty$ 

- So M = 0 must be a maximum! - There are no max/min's to the left of M= 0, and f(M) is decreasing to  $-\infty$ . - So  $M = \frac{4}{5}E_{\text{tot}}$  must be a minimum!

- There are no max/mins to the right of M

 $=\frac{4}{5}E_{\text{total}}$ , and f(M) is increasing to  $+\infty$ .



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 $=\frac{4}{5}E_{\text{total}}$ , and f(M) is increasing to  $+\infty$ .



- Let  $f(M) = M^4(M E_{\text{total}}) + \beta V$   $f(M) = 0 \Leftrightarrow T_{\text{BH}} = T_{\text{gas}}$ 
  - <u>So</u>: Condition for there to be physically reasonable values of *M* for which  $T_{BH} = T_{gas}$ :

$$f(\frac{4}{5}E_{\text{total}}) \le 0$$
 or  $\beta V \le 0.082E_{\text{tot}}^5$ 

Are these values of M stable equilibrium points?  $\frac{dT_{\rm BH}}{dM} = \frac{d}{dM} \left(\frac{1}{8\pi M}\right) = -\frac{1}{8\pi M^2}$ Do they satisfy Condition (2)  $\frac{dT_{\rm BH}}{dM} = \frac{d}{dM} \left(\frac{1}{8\pi M}\right) = -\frac{1}{8\pi M^2}$ 

$$\frac{dT_{gas}}{dM} = \frac{d}{dM} \left( \left[ \frac{E_{\text{total}} - M}{\frac{\pi^2}{15}V} \right]^{1/4} \right) = -\frac{1}{4} \left( \frac{\pi^2}{15}V \right)^{-1/4} (E_{\text{total}} - M)^{-3/4}$$

So Condition (2) is: 
$$-\frac{1}{8\pi M^2} < -\frac{1}{4} \left(\frac{\pi^2}{15}V\right)^{-1/4} (E_{\text{total}} - M)^{-3/4}$$

$$\underline{Or}: \quad \frac{1}{4^4} \frac{M^8}{(E_{\text{total}} - M)^3} < \beta V$$

- For physically reasonable values  $\overline{4}$ of *M*, this becomes:

$$\frac{1}{4^4} \frac{M^8}{(E_{\text{total}} - M)^3} \le 0.082 E_{\text{total}}^5$$

or 
$$M < \frac{4}{5}E_{\text{total}}$$

- <u>Case 1</u>:  $f(\frac{4}{5}E_{\text{total}}) < 0, 0 < \beta V < 0.082E_{\text{tot}}^{5}$ 



# **3. A Black Hole Carnot Cycle**

- <u>So</u>: It looks like a black hole *can* be in equilibrium with a photon gas, and hence potentially undergo a reversible process.
- *Next task*: Describe such a process.
  - Construct a Carnot cycle using a black hole and photon gas as the working fluid.
- *Slight catch*: Black hole + photon gas working fluid has negative heat capacity.

To use it to power a heat engine, need to reverse the standard (positive heat capacity working fluid) Carnot cycle.







A reverse positive heat capacity heat engine!

### <u>Negative heat capacity Carnot cycle heat engine</u>

- heat absorbed  $\rightarrow$  temp decrease
- heat radiated  $\rightarrow$  temp increase

Heat engine: takes heat from a hot place to a cold place and produces work



A "reverse" positive heat capacity refrigerator!



#### <u>Tasks</u>:

(a) Show that efficiency  $\mu = (Q_{\rm in} - Q_{\rm out})/Q_{\rm in} = 1 - T_2/T_1$ 

(b) Show that  $S_{\text{TD}}$  of BH subsystem is  $1/16\pi T^2 = 4\pi M^2 = A/4 = S_{\text{BH}}$ 

$$E_{\rm BH+gas} = U_{\rm BH+gas} = M + \alpha V T^4, \quad \alpha = \frac{\pi^2}{15}$$

$$S_{\rm gas} = \int dU_{\rm gas} / T = \int (4\alpha V T^3 dT) / T = \int 4\alpha V T^2 dT = \frac{4}{3} \alpha V T^3$$

$$P_{\rm gas} = \frac{1}{3} (U_{\rm gas} / V) = \frac{1}{3} \alpha T^4$$

 $(1 \rightarrow 2) Isothermal Expansion$  $V_1 \rightarrow V_2 at constant T = T_1$ 

• Note: 
$$\delta Q_{12} = dU_{BH+gas} + P_{gas}dV$$
$$= (\partial U_{BH+gas}/\partial V)_T dV + \frac{1}{3}\alpha T_1^4 dV$$
$$= \alpha T_1^4 dV + \frac{1}{3}\alpha T_1^4 dV$$
$$\bullet \underline{So}: \qquad Q_{12} = \int_{V_1}^{V_2} (\alpha T_1^4 + \frac{1}{3}\alpha T_1^4) dV = \frac{4}{3}\alpha T_1^4 (V_2 - V_1)$$
$$\bullet \underline{And}: \qquad \Delta S_{12} = \int_{1}^{2} (\delta Q_{12}/T_1) = \frac{4}{3}\alpha T_1^3 (V_2 - V_1)$$
$$\bullet \underline{C} = Resulting change in entropy$$

 $T_1$ 

Q

2

4

 $T_2$ 

3

 $Q_{12}$ 

$$E_{\rm BH+gas} = U_{\rm BH+gas} = M + \alpha V T^4, \quad \alpha = \frac{\pi^2}{15}$$

$$S_{\rm gas} = \int dU_{\rm gas} / T = \int (4\alpha V T^3 dT) / T = \int 4\alpha V T^2 dT = \frac{4}{3} \alpha V T^3$$

$$P_{\rm gas} = \frac{1}{3} (U_{\rm gas} / V) = \frac{1}{3} \alpha T^4$$

$$(2 \rightarrow 3) A diabatic CompressionQ_{23} = 0, V_2 \rightarrow V_3, T_1 \rightarrow T_2$$

• So: 
$$\Delta S_{23} = \int_{2}^{3} (\delta Q_{23}/T_1) = 0$$

• <u>Adiabatic relation</u>:  $dU_{\rm BH+gas} = dW_{\rm gas}$  $dM + 4\alpha VT^3 dT + \alpha T^4 dV$   $-\frac{1}{3}\alpha T^4 dV =$  work done on gas

$$Or... \quad VT^3 + \frac{3\gamma}{8\alpha T^2} = \text{const.}$$

-  $\gamma = \frac{1}{8\pi}$ - For photon gas by itself, adiabatic relation is just  $VT^3 = \text{const.}$ 

Or... 
$$V_3 T_2^3 + \frac{3\gamma}{8\alpha T_2^2} = V_2 T_1^3 + \frac{3\gamma}{8\alpha T_1^2}$$



$$E_{\rm BH+gas} = U_{\rm BH+gas} = M + \alpha V T^4, \quad \alpha = \frac{\pi^2}{15}$$
  

$$S_{\rm gas} = \int dU_{\rm gas} / T = \int (4\alpha V T^3 dT) / T = \int 4\alpha V T^2 dT = \frac{4}{3} \alpha V T^3$$
  

$$P_{\rm gas} = \frac{1}{3} (U_{\rm gas} / V) = \frac{1}{3} \alpha T^4$$

 $(3 \rightarrow 4) Isothermal Compression$  $V_3 \rightarrow V_4 at constant T = T_2$ 

• In derviation similar to  $(1 \rightarrow 2)$ :

$$\frac{1}{Q_{34}} = \int_{V_3}^{V_4} \left(\alpha T_2^4 + \frac{1}{3}\alpha T_2^4\right) dV = \frac{4}{3}\alpha T_2^4 (V_4 - V_3)$$

$$AS_{34} = \int_3^4 \left(\delta Q_{34}/T_2\right) = \frac{4}{3}\alpha T_2^3 (V_4 - V_3)$$

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$$C = Resulting change in entropy$$

 $T_1$ 

Q

2

 $Q_{12}$ 

$$E_{\rm BH+gas} = U_{\rm BH+gas} = M + \alpha V T^4, \quad \alpha = \frac{\pi^2}{15}$$

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$$P_{\rm gas} = \frac{1}{3} (U_{\rm gas} / V) = \frac{1}{3} \alpha T^4$$

 $\frac{(4 \rightarrow 1) A diabatic Expansion}{Q_{41} = 0, V_4 \rightarrow V_1, T_2 \rightarrow T_1}$ 

• So: 
$$\Delta S_{41} = \int_{4}^{1} (\delta Q_{41}/T_2) = 0$$

• <u>Recall adiabatic relation for  $(2 \rightarrow 3)$ </u>:  $V_3 T_2^3 + \frac{3\gamma}{8\alpha T_2^2} = V_2 T_1^3 + \frac{3\gamma}{8\alpha T_1^2}$ • <u>For  $(4 \rightarrow 1)$  this becomes</u>:  $V_4 T_2^3 + \frac{3\gamma}{8\alpha T_2^2} = V_1 T_1^3 + \frac{3\gamma}{8\alpha T_1^2}$  $V_4 T_2^3 - V_3 T_2^3 = V_1 T_1^3 - V_2 T_1^3$ 



Task (a): Show that efficiency  $\mu = (Q_{in} - Q_{out})/Q_{in} = 1 - T_2/T_1$ 

$$\mu = \frac{Q_{12} - Q_{34}}{Q_{12}}$$

$$= \frac{\frac{4}{3}\alpha T_1^4 (V_2 - V_1) - \frac{4}{3}\alpha T_2^4 (V_4 - V_3)}{Q_{12}}$$

$$= \frac{\frac{4}{3}\alpha [T_1 (V_2 T_1^3 - V_1 T_1^3)] - \frac{4}{3}\alpha [T_2 (V_4 T_2^3 - V_3 T_2^3)]}{Q_{12}}$$

$$= \frac{\frac{4}{3}\alpha [T_1 (V_2 T_1^3 - V_1 T_1^3)] - \frac{4}{3}\alpha [T_2 (V_1 T_1^3 - V_2 T_1^3)]}{Q_{12}}$$

$$= \frac{\frac{4}{3}\alpha T_1^4 (V_2 - V_1) \left(1 - \frac{T_2}{T_1}\right)}{Q_{12}}$$

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$$= \frac{4}{3}\alpha T_1^4 (V_2 - V_1) \left(1 - \frac{T_2}{T_1}\right)$$

 $= 1 - \frac{T_2}{T_1}$  Task (a) accomplished!

<u>Task (b): Show that  $S_{TD}$  of BH subsystem is  $1/16\pi T^2$ </u>

• Isothermal processes  $(1\rightarrow 2)$ ,  $(3\rightarrow 4)$ :

 $\Delta S_{12} = \frac{4}{3} \alpha T_1^{3} (V_2 - V_1)$   $\Delta S_{34} = \frac{4}{3} \alpha T_2^{3} (V_4 - V_3)$ Entropy change in gas only (no change of state of BH)

- Adiabatic processes (2 $\rightarrow$ 3), (4 $\rightarrow$ 1):  $\Delta S_{BH+gas} = 0$ 
  - <u>But</u>: Adiabatic relation of gas by itself is different from adiabatic relation of BH+gas.
  - <u>Claim</u>: There must be heat flux from gas to BH during adiabatic compression  $(2\rightarrow 3)$ .

• So: 
$$\Delta S_{BH+gas,23} = \Delta S_{BH,23} + \Delta S_{gas,23}$$
  

$$= \Delta S_{BH,23} + \frac{4}{3}\alpha (V_3 T_2{}^3 - V_2 T_1{}^3)$$

$$= \Delta S_{BH,23} + \frac{4}{3}\alpha \left(\frac{3\gamma}{8\alpha T_2{}^2} - \frac{3\gamma}{8\alpha T_1{}^2}\right)$$

$$= \Delta S_{BH,23} + \frac{1}{2}\gamma \left(\frac{1}{T_2{}^2} - \frac{1}{T_1{}^2}\right) = 0$$

$$V_3 T_2{}^3 + \frac{3\gamma}{8\alpha T_2{}^2} = V_2 T_1{}^3 + \frac{3\gamma}{8\alpha T_1{}^2}$$

- <u>Or</u>:  $\Delta S_{\rm BH,23} = \frac{1}{2} \gamma \left( \frac{1}{T_1^2} \frac{1}{T_2^2} \right)$
- <u>Ingeneral</u>:  $S_{\rm BH} = \frac{1}{2}\gamma(1/T^2) = 1/16\pi T^2$

Task (b) accomplished!

 $T_1$ 

0

2

3

 $Q_{12}$ 

Q

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