

# 08. Quantum Mechanics: Density Operators

1. States as Density Operators
2. Pure States and Mixed States
3. Test for Mixedness

## 1. States as Density Operators

- Motivation: Suppose a quantum system is associated with an *ensemble*  $\{|\psi_i\rangle, p_i\}$  of vector states  $|\psi_i\rangle$ , each with a probability  $p_i$ .



*Suppose we don't know what state our quantum system is in; suppose all we know are its possible states.*

- Recall: In Gibbs' approach to classical statistical mechanics, we considered an ensemble of *classical* states to be a *classical* phase space  $\Gamma$  with a probability distribution  $\rho(x, t)$  defined on it.

- *In QM, we use a vector space  $\mathcal{H}$  instead of a classical phase space  $\Gamma$ .*
- *What should we use instead of the function  $\rho(x, t)$ ?*

*Answer: An operator...*

**Def. 1** (*Density operator*). The **density operator**  $\rho$  for a system in an ensemble  $\{|\psi_i\rangle, p_i\}$  of  $m$  vector states  $|\psi_i\rangle \in \mathcal{H}$  each with probability  $p_i$ , is defined by

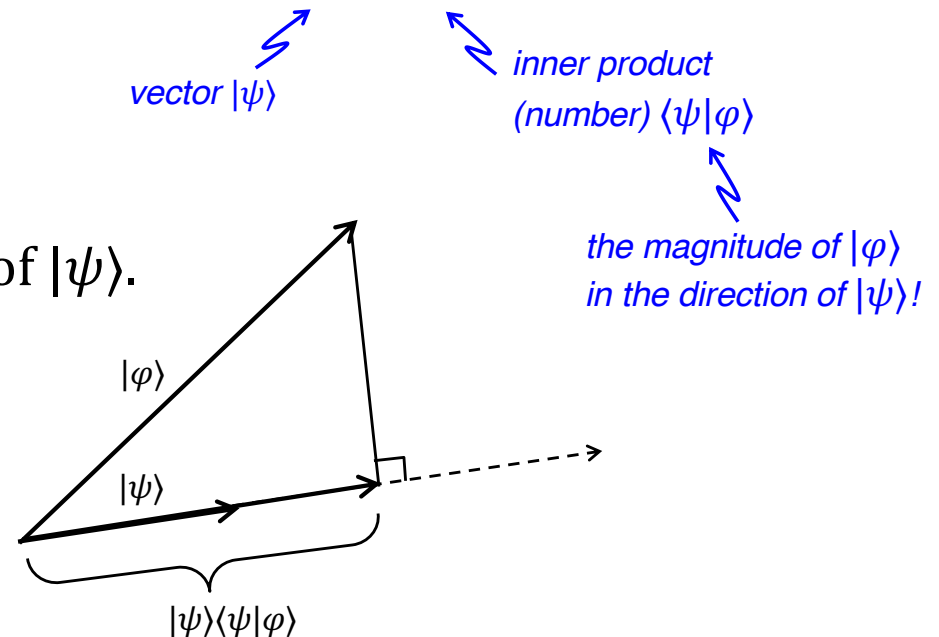
$$\rho = \sum_{i=1}^m p_i |\psi_i\rangle \langle \psi_i|, \text{ where } \sum_{i=1}^m p_i = 1.$$

Warning! Same notation but different mathematical objects:  
 $\rho$  in QM is an operator.  
 $\rho(x, t)$  in classical Stat Mech is a function.

- Note 1:  $|\psi\rangle \langle \psi|$  is an operator.
  - It acts on a vector  $|\varphi\rangle$  and gives you another vector  $|\psi\rangle \langle \psi| \varphi\rangle$ .

- Note 2:  $|\psi\rangle \langle \psi|$  is a "projection" operator.

- It projects a vector  $|\varphi\rangle$  in the direction of  $|\psi\rangle$ .



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- Note 3: The vector states  $\{|\psi_i\rangle\}$  are not necessarily orthogonal to each other, and  $m$  is not necessarily the dimension  $n$  of  $\mathcal{H}$ .

But:  $\rho$  can always be expressed in terms of a basis  $\{|\phi_i\rangle\}$  of  $\mathcal{H}$  as

$$\rho = \sum_{i=1}^n \lambda_i |\phi_i\rangle \langle \phi_i|$$

where  $|\phi_i\rangle$  are eigenvectors of  $\rho$  with eigenvalues  $\lambda_i \geq 0$ .

← Because  $\rho$  is Hermitian!

- And: Since  $\rho$  is supposed to represent an ensemble of vector states, in this case  $\{|\phi_i\rangle, \lambda_i\}$ , we require  $\sum_{i=1}^n \lambda_i = 1$ . ← The sum of the eigenvalues of a density operator is 1!

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- Note 4: A density operator can describe more than one ensemble of vector states.

Example:

$$\begin{aligned} \text{Suppose: } |a\rangle &= \sqrt{3/4}|0\rangle + \sqrt{1/4}|1\rangle & \text{Then: } |0\rangle &= \sqrt{1/3}|a\rangle + \sqrt{1/3}|b\rangle \\ |b\rangle &= \sqrt{3/4}|0\rangle - \sqrt{1/4}|1\rangle & |1\rangle &= |a\rangle - |b\rangle \end{aligned}$$

$$\text{So: } \rho = \frac{1}{2}|a\rangle\langle a| + \frac{1}{2}|b\rangle\langle b| \quad \text{corresponds to } \rightsquigarrow \{|a\rangle, |b\rangle; \frac{1}{2}, \frac{1}{2}\}$$

$$= \frac{3}{4}|0\rangle\langle 0| + \frac{1}{4}|1\rangle\langle 1| \quad \text{corresponds to } \rightsquigarrow \{|0\rangle, |1\rangle; \frac{3}{4}, \frac{1}{4}\}$$

↖ distinct ensembles,  
same density operator

## 2. Pure States and Mixed States

*Why use density operators to represent states?*

**Def. 2** (*Pure/mixed density operator state*). If a density operator can be expressed as  $\rho = |\psi\rangle\langle\psi|$ , then it is called a **pure density operator state**. Otherwise, it is called a **mixed density operator state**.

- Note 1: An "ignorance" interpretation of mixed density operator states is not always appropriate!

$$\rho = p_1|\psi_1\rangle\langle\psi_1| + p_2|\psi_2\rangle\langle\psi_2| + \dots$$


 *Ignorance as to what vector state the system is in?*

- Concern 1:  $\rho$  can describe more than one ensemble of vector states.

Recall:

Let:  $|a\rangle = \sqrt{3/4}|0\rangle + \sqrt{1/4}|1\rangle$

$$|b\rangle = \sqrt{3/4}|0\rangle - \sqrt{1/4}|1\rangle$$

 *If the system is in the density operator state  $\rho$ , should it be associated with the ensemble  $\{|a\rangle, |b\rangle; 1/2, 1/2\}$  or the ensemble  $\{|0\rangle, |1\rangle; 3/4, 1/4\}$ ?*

Then:  $\rho = 1/2|a\rangle\langle a| + 1/2|b\rangle\langle b|$  corresponds to  $\{|a\rangle, |b\rangle; 1/2, 1/2\}$

$$= 3/4|0\rangle\langle 0| + 1/4|1\rangle\langle 1| \text{ corresponds to } \{|0\rangle, |1\rangle; 3/4, 1/4\}$$

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 *Ignorance as to what vector state the system is in?*

- Concern 2: An ignorance interpretation cannot be applied to a mixed density operator state of a subsystem of a composite system in a pure *entangled* density operator state.

 *To be demonstrated later...*

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**Def. 2** (*Pure/mixed density operator state*). If a density operator can be expressed as  $\rho = |\psi\rangle\langle\psi|$ , then it is called a **pure density operator state**. Otherwise, it is called a **mixed density operator state**.

- Note 2: A mixed density operator state is not the same thing as a superposed vector state!


Example:

- The mixed density operator state  $\rho = \frac{1}{2}|a\rangle\langle a| + \frac{1}{2}|b\rangle\langle b|$  corresponds to an *ensemble* of vector states  $\{|a\rangle, |b\rangle; \frac{1}{2}, \frac{1}{2}\}$ .
- The pure density operator state  $\rho = \frac{1}{2}\{|a\rangle + |b\rangle\}\{\langle a| + \langle b|\}$  corresponds to a *single* vector state  $\frac{1}{\sqrt{2}}\{|a\rangle + |b\rangle\}$ .

### 3. A Test for Mixedness


**Def. 3 (Trace).** Let  $O$  be an operator on a vector space  $\mathcal{H}$  with basis  $\{|\nu_i\rangle\}$ . The **trace**  $\text{Tr}(O)$  of  $O$  is defined by

$$\text{Tr}(O) \equiv \sum_i \langle \nu_i | O | \nu_i \rangle$$

 *The sum of the diagonal elements of a matrix representation of  $O$*

- Note 1: If  $O$  is a Hermitian operator on  $\mathcal{H}$ , then its eigenvectors  $\{|\lambda_i\rangle\}$  form a basis for  $\mathcal{H}$ , and

$$\text{Tr} O = \sum_i \langle \lambda_i | O | \lambda_i \rangle = \sum_i \lambda_i \langle \lambda_i | \lambda_i \rangle = \sum_i \lambda_i$$

 *The trace of a Hermitian operator is the sum of its eigenvalues!*

- Note 2: Recall that the sum of the eigenvalues of a density operator state  $\rho$  is 1.
  - So:  $\text{Tr} \rho = 1$  for both pure and mixed density operator states.



**Claim** (*Test for mixedness*).

- (a)  $\rho$  is a pure density operator state *if and only if*  $\text{Tr } \rho^2 = 1$ .  
 (b)  $\rho$  is a mixed density operator state *if and only if*  $\text{Tr } \rho^2 < 1$ .

• First note:

$$\begin{aligned}
 \rho^2 &= \sum_{i=1}^n \lambda_i |\phi_i\rangle\langle\phi_i| \sum_{j=1}^n \lambda_j |\phi_j\rangle\langle\phi_j| \quad \leftarrow \sum |\phi_i\rangle \text{ are eigenvectors of } \rho \text{ with eigenvalues } \lambda_i \\
 &= \sum_{i,j} \lambda_i \lambda_j |\phi_i\rangle\langle\phi_i|\phi_j\rangle\langle\phi_j| \quad \leftarrow \sum \langle\phi_i|\phi_j\rangle = 0 \text{ unless } j = i \\
 &= \sum_i \lambda_i^2 |\phi_i\rangle\langle\phi_i| \quad \leftarrow \sum |\phi_i\rangle \text{ are eigenvectors of } \rho^2 \text{ with eigenvalues } \lambda_i^2 \quad \searrow \text{Tr } \rho^2 = \sum_i \lambda_i^2
 \end{aligned}$$

Proof of (a):

- Suppose:  $\rho$  is a pure density operator state.
  - Then:  $\rho^2 = \rho$ , and hence  $\text{Tr } \rho^2 = \text{Tr } \rho = 1$ .
- Now suppose:  $\rho$  is a density operator state and  $\text{Tr } \rho^2 = 1$ .
  - Then:  $\sum_i \lambda_i^2 = 1 = \sum_i \lambda_i$ .
  - And: This holds if and only if one of the  $\lambda_i$  is 1 and the rest are 0.
  - Thus:  $\rho$  is pure.

**Claim** (*Test for mixedness*).

- (a)  $\rho$  is a pure density operator state *if and only if*  $\text{Tr } \rho^2 = 1$ .
- (b)  $\rho$  is a mixed density operator state *if and only if*  $\text{Tr } \rho^2 < 1$ .

• First note:

$$\begin{aligned}
 \rho^2 &= \sum_{i=1}^n \lambda_i |\phi_i\rangle \langle \phi_i| \sum_{j=1}^n \lambda_j |\phi_j\rangle \langle \phi_j| \quad \leftarrow \sum |\phi_i\rangle \text{ are eigenvectors of } \rho \text{ with eigenvalues } \lambda_i \\
 &= \sum_{i,j} \lambda_i \lambda_j |\phi_i\rangle \langle \phi_i | \phi_j \rangle \langle \phi_j| \quad \leftarrow \sum \langle \phi_i | \phi_j \rangle = 0 \text{ unless } j = i \\
 &= \sum_i \lambda_i^2 |\phi_i\rangle \langle \phi_i| \quad \leftarrow \sum |\phi_i\rangle \text{ are eigenvectors of } \rho^2 \text{ with eigenvalues } \lambda_i^2 \quad \searrow \text{Tr } \rho^2 = \sum_i \lambda_i^2
 \end{aligned}$$

Proof of (b):

- Suppose:  $\rho$  is a mixed density operator state.
  - Then:  $\text{Tr } \rho^2 = \sum_i \lambda_i^2 \leq \sum_i \lambda_i = 1$ , with equality if and only if  $\rho$  is pure.
- Now suppose:  $\rho$  is a density operator state and  $\text{Tr } \rho^2 < 1$ .
  - Then:  $\text{Tr } \rho^2 = \sum_i \lambda_i^2 < 1$ , and  $\sum_i \lambda_i = 1$ .
  - And: This excludes the case of one of the  $\lambda_i$  being 1 and the rest 0.
  - Thus:  $\rho$  must be mixed.

*Why get bogged down with all these linear algebra definitions of mixed density operator states and tests of mixedness?*

- Because:

- Ultimately: We're going to focus on a notion of entropy associated with density operator states (think of the analogy with the Gibbs entropy for classical ensembles).
- And: That notion of entropy (called the "von Neumann" entropy) is first and foremost a measure of the degree to which a density operator state is mixed!
- Moreover: As we'll see, it can also be interpreted (under certain conditions) as a measure of *entanglement*.

  
*Next up...*