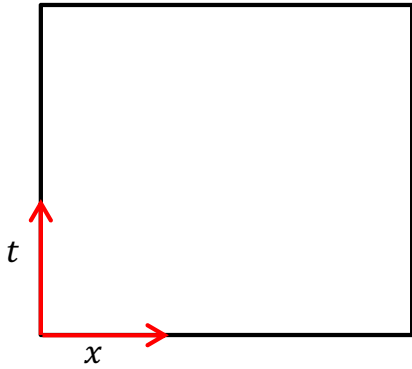


# 11. Spacetime

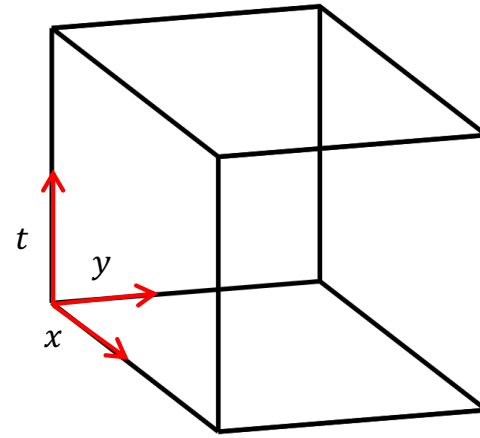
1. Types of Spacetime
2. Classical Spacetimes

## 1. Types of Spacetimes

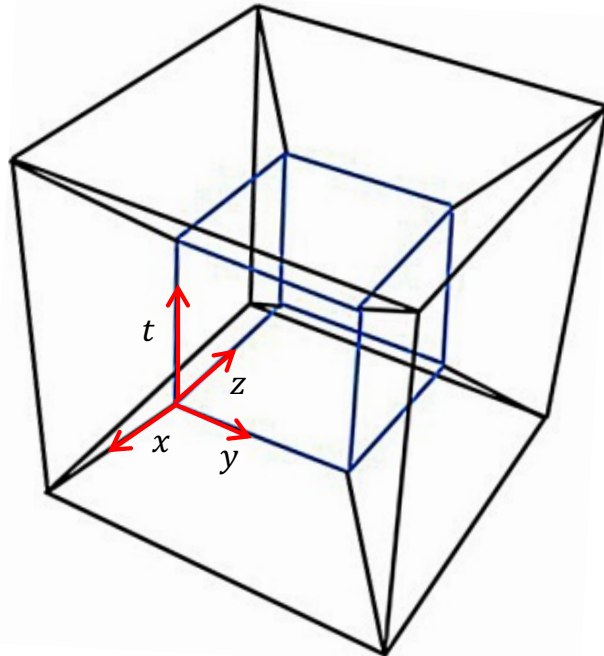
- A *spacetime* is a 4-dim collection of points with additional structure.



2-dim square



"shadow" (projection)  
of 3-dim cube in 2-dim



"shadow" (projection)  
of 4-dim hyper-cube  
("tesseract") in 2-dim

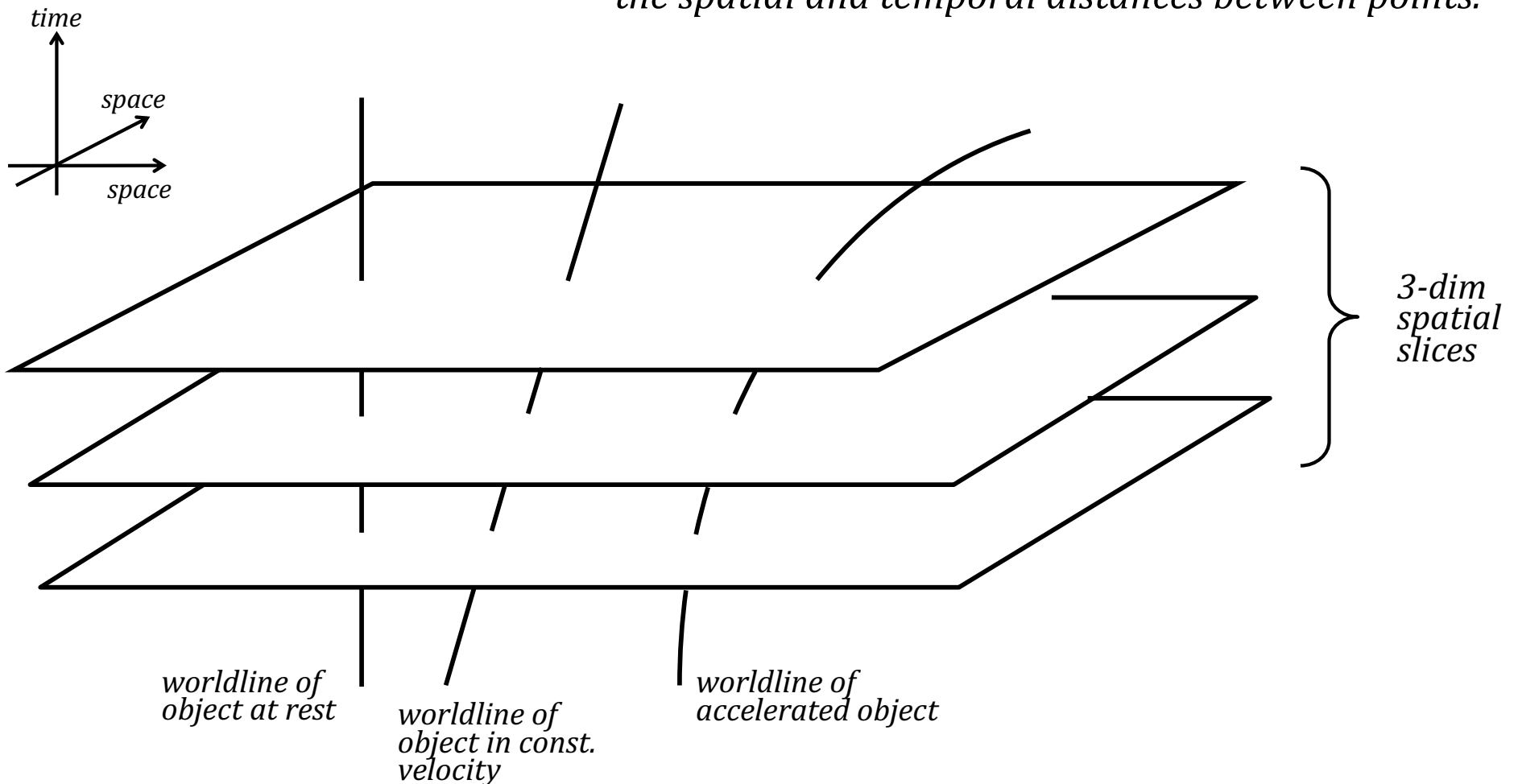
# 11. Spacetime

1. Types of Spacetime
2. Classical Spacetimes

## 1. Types of Spacetimes

- A *spacetime* is a 4-dim collection of points with additional structure.

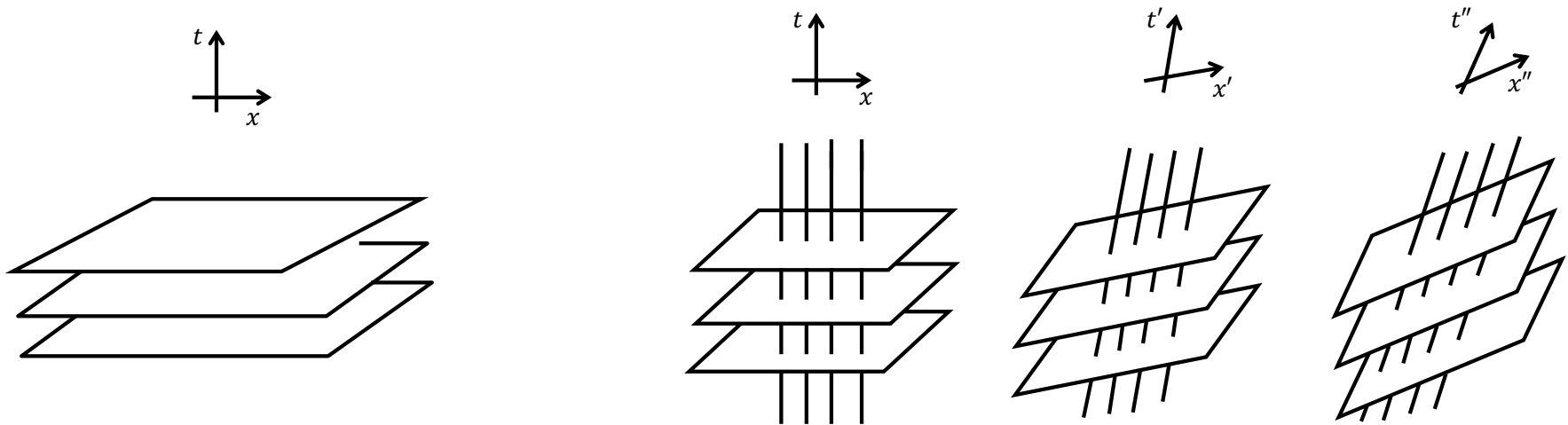
Typically, one or more metrics = a specification of the spatial and temporal distances between points.



Two ways spacetimes can differ:

(1) Different ways of specifying distances between points yield different types of spacetimes.

- *Classical spacetimes* have *separate* spatial and temporal metrics: only one way to split time from space (spatial and temporal distances are *absolute*).
- *Relativistic spacetimes* have a *single* spatiotemporal metric, and how it gets split into spatial and temporal parts depends on one's inertial reference frame (spatial and temporal distances are *relative*).



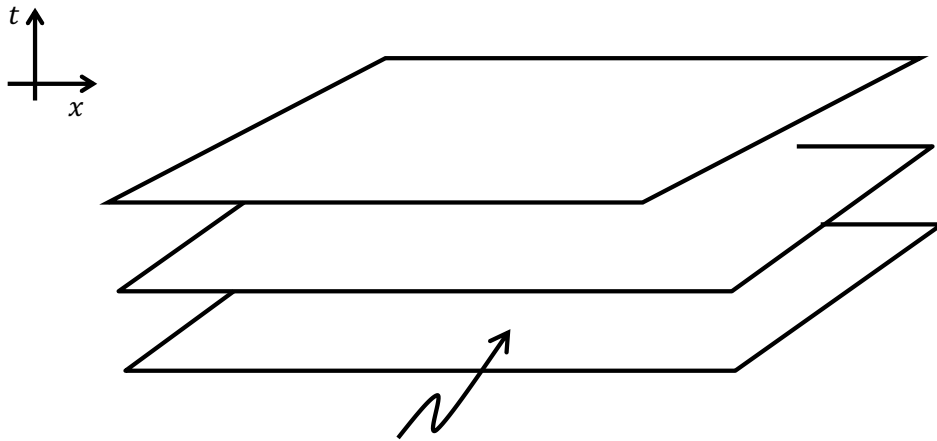
*Classical spacetime: only one way to split time from space.*

*Relativistic spacetime: many ways to split time from space.*

(2) Metrics can be *flat* or *curved*: how one specifies the distance between points encodes the curvature of the spacetime.

- *Classical spacetimes* can be flat or curved.
- *Relativistic spacetimes* can be flat (Minkowski spacetime) or curved (general relativistic spacetimes).

Two ways curvature can manifest itself



*How the spatial slices are "rigged" together can be flat or curved.*

*The spatial slices can be flat or curved.*

## 2. Classical Spacetimes

Newtonian spacetime is a 4-dim collection of points such that:

(N1) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *temporal interval*  $T(p, q) = t' - t$ .

(N2) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *Euclidean distance*

$$R(p, q) = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' - z)^2}$$

(N1) and (N2) entail:

(a) All worldlines have a definite *absolute velocity*.

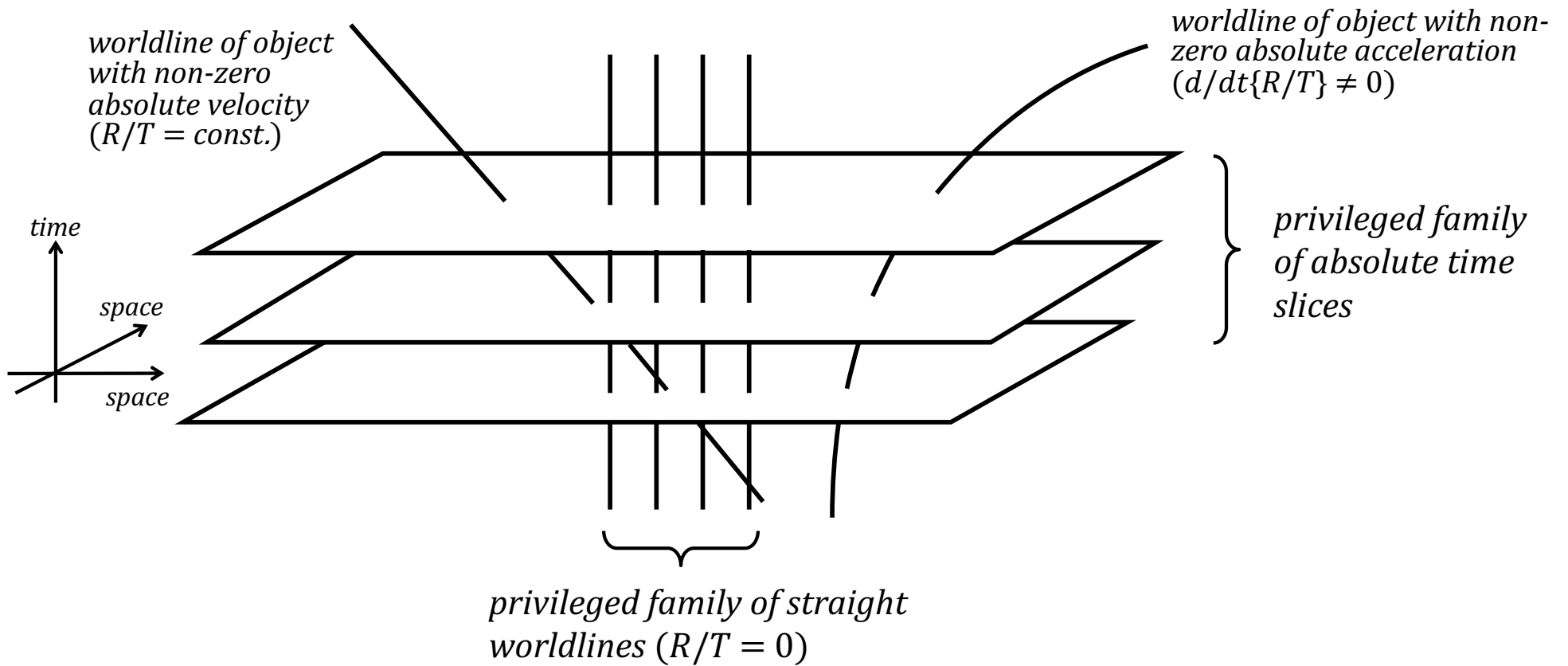
For worldline  $\gamma$ , and any two points  $p, q$  on  $\gamma$ , the *absolute velocity* of  $\gamma$  with respect to  $p, q$  can be defined by  $R(p, q)/T(p, q)$ .

(b) There is a privileged collection of worldlines defined by  $R(p, q)/T(p, q) = 0$ .

(c) All worldlines have a definite *absolute acceleration*.

For worldline  $\gamma$ , and points  $p, q$  on  $\gamma$ , the *absolute acceleration* of  $\gamma$  with respect to  $p, q$  can be defined by  $d/dt\{R(p, q)/T(p, q)\}$ .

# Newtonian Spacetime



1. Single, privileged inertial frame.
2. Velocity is absolute.
3. Acceleration is absolute.
4. Simultaneity is absolute.

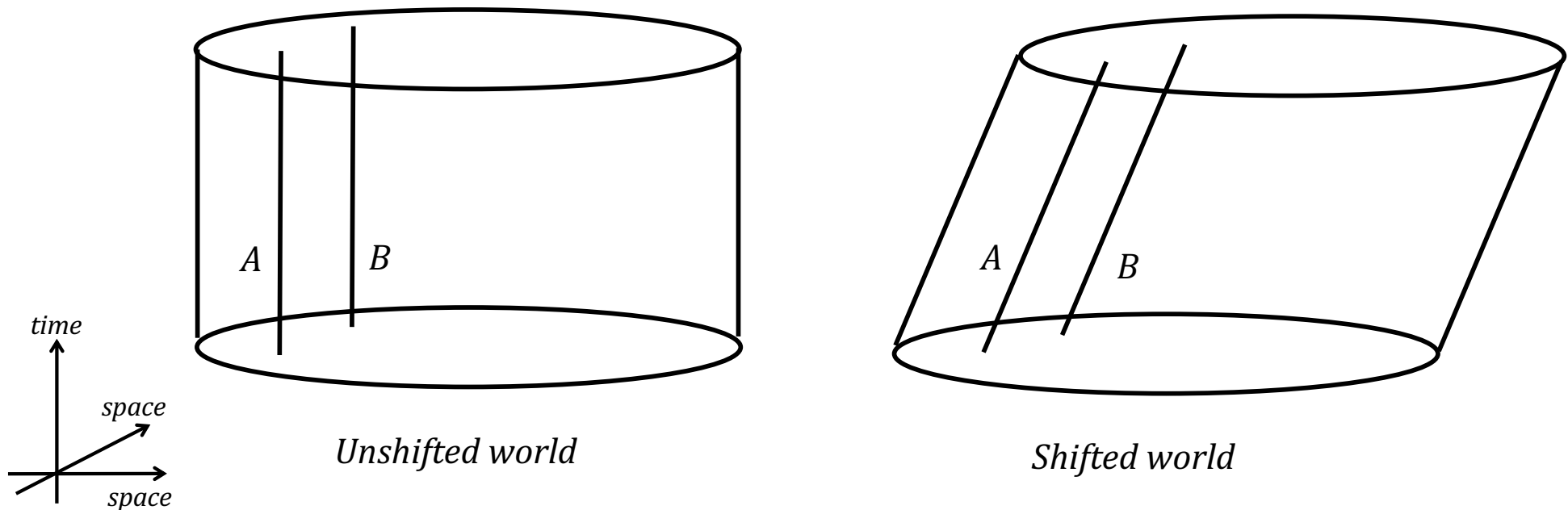
- Newtonian spacetime has enough structure to support the Principle of Inertia.

*Principle of Inertia in Newtonian Spacetime*

Objects follow straight worldlines ( $R/T = \text{const}$ ) unless acted upon by external forces.

- *But*: Newtonian spacetime also supports Leibiz's Kinematic Shift!

*Kinematic Shift in Newtonian Spacetime*



- In Newtonian spacetime, the unshifted and shifted worlds are distinct!
  - Families of straight lines with different slopes can be distinguished from each other.

Galilean spacetime is a 4-dim collection of points such that:

(G1) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *temporal interval*  $T(p, q) = t' - t$ .

(G2) Between any two *simultaneous* points  $p_t, q_t$ , with coordinates  $(t, x, y, z)$  and  $(t, x', y', z')$ , there is a definite *Euclidean distance*,

$$R(p_t, q_t) = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' - z)^2}$$

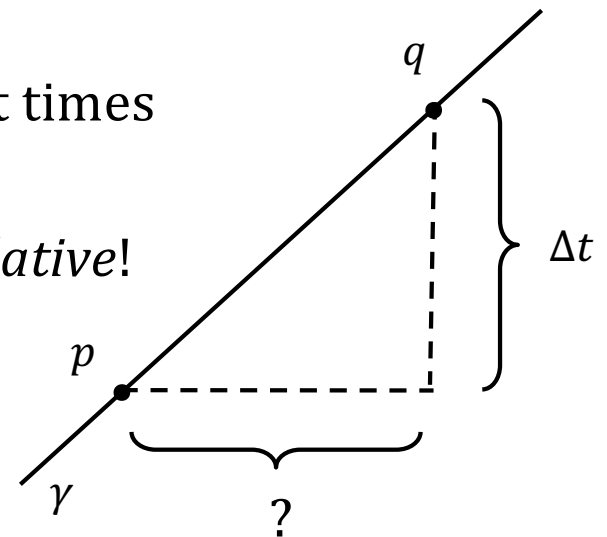
(G3) Any worldline  $\gamma$  through point  $p$  has a definite curvature  $S(\gamma, p)$ .

(G2) entails:

- No absolute spatial distance between points at different times on any worldline  $\gamma$ .
- So: No absolute velocity for any worldline: *velocity is relative!*
- So: No *single* privileged frame of reference.

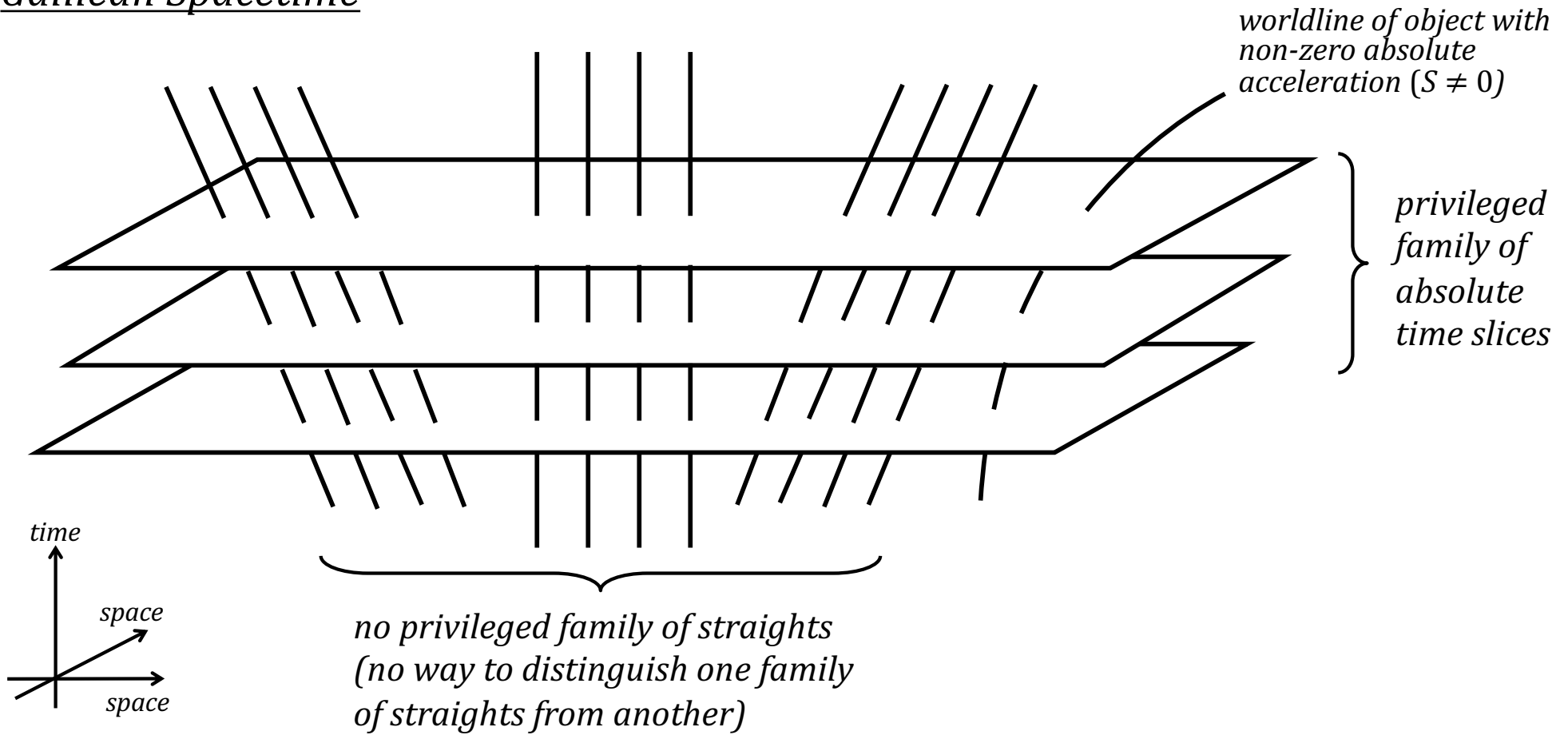
(G3) entails: acceleration remains absolute!

For worldline  $\gamma$  and point  $p$  on  $\gamma$ , the *absolute acceleration* of  $\gamma$  with respect to  $p$  is given by  $S(\gamma, p)$ .





# Galilean Spacetime



1. Many inertial frames; none privileged.
2. Velocity is relative.
3. Acceleration is absolute.
4. Simultaneity is absolute.

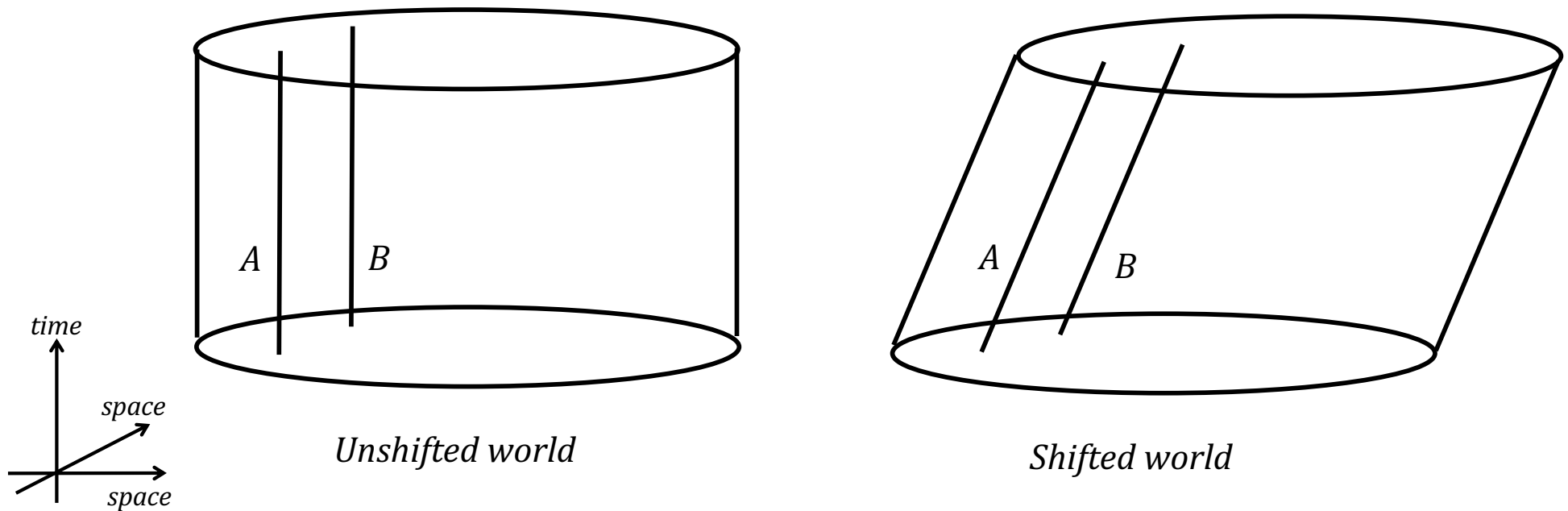
- Galilean spacetime has enough structure to support the Principle of Inertia.

*Principle of Inertia in Galilean Spacetime*

Objects follow straight worldlines ( $S = 0$ ) unless acted upon by external forces.

- And: Galilean spacetime does *not* support Leibiz's Kinematic Shift!

*Kinematic Shift in Galilean Spacetime*

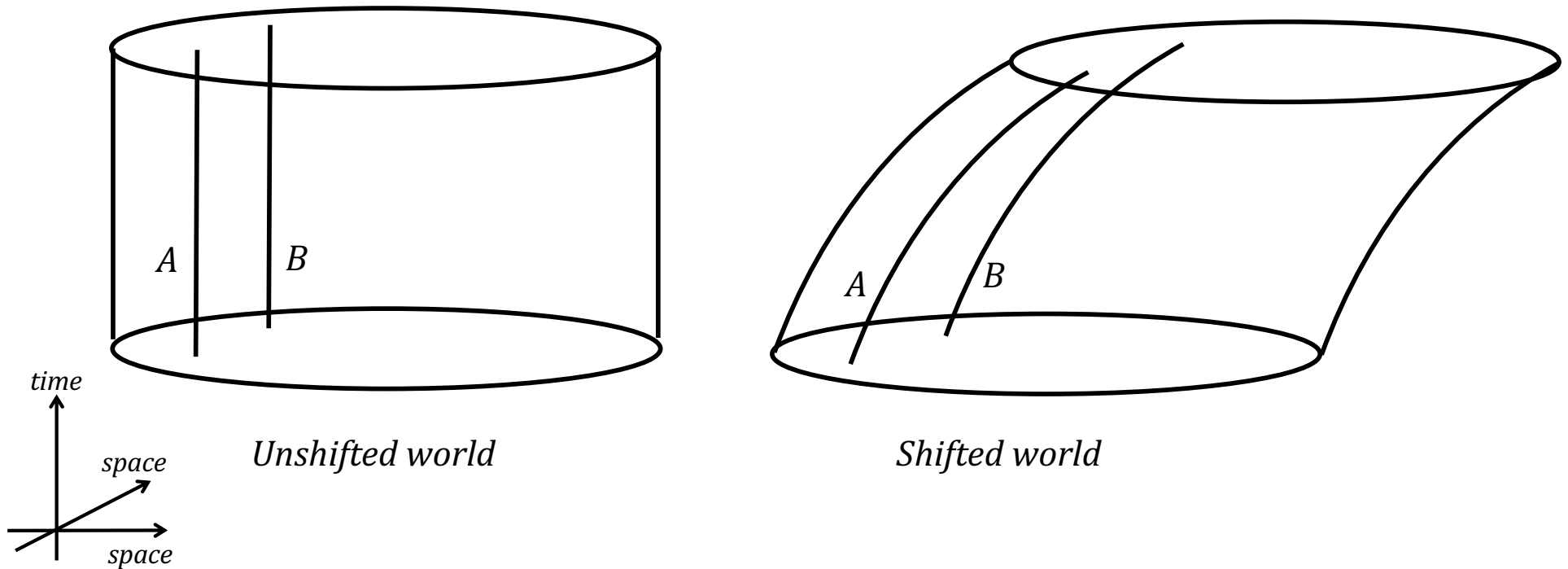


- In Galilean spacetime, the unshifted and shifted worlds are indistinguishable! (Families of straight lines cannot be distinguished from each other.)

## What about Clarke's Dynamic Shift?

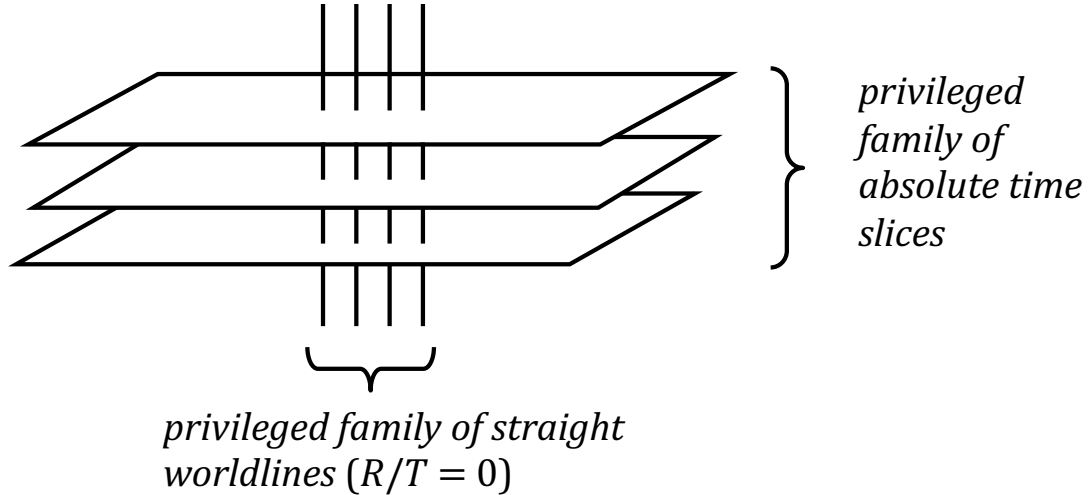
- Both Newtonian and Galilean spacetimes support the Dynamic Shift:

### Dynamic Shift in Newtonian and Galilean Spacetimes



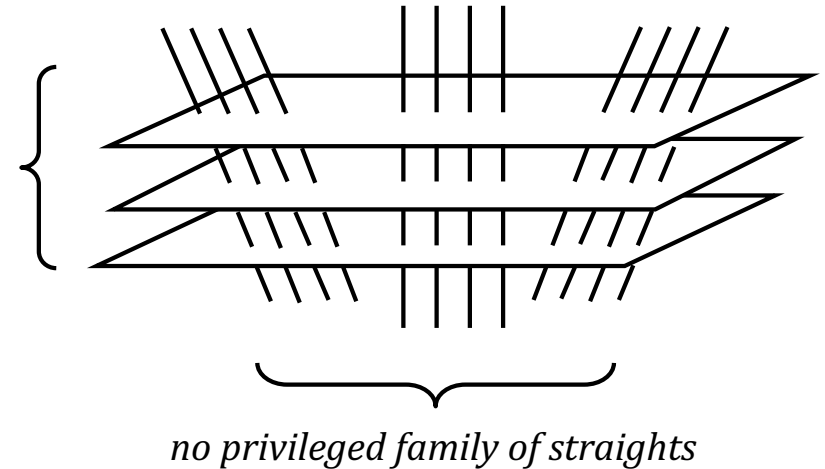
- In Newtonian spacetime, the unshifted and shifted worlds differ on their values of absolute acceleration  $d/dt\{R/T\}$ .
- In Galilean spacetime, the unshifted and shifted worlds differ on their values of absolute acceleration  $S$ .
- (In both Newtonian and Galilean spacetime, straight worldlines are distinct from curved worldlines.)

## Newtonian Spacetime



1. Single, privileged inertial frame.
2. Velocity is absolute.
3. Acceleration is absolute.
4. Simultaneity is absolute.

## Galilean Spacetime



1. Many inertial frames; none privileged.
2. Velocity is relative.
3. Acceleration is absolute.
4. Simultaneity is absolute.

- There are no privileged locations in Newtonian and Galilean spacetimes (they are homogeneous).
- Recall: Aristotle's cosmos has a privileged location...

Aristotelian spacetime is a 4-dim collection of points such that:

(A1) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *temporal interval*  $T(p, q) = t' - t$ .

(A2) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *Euclidean distance*

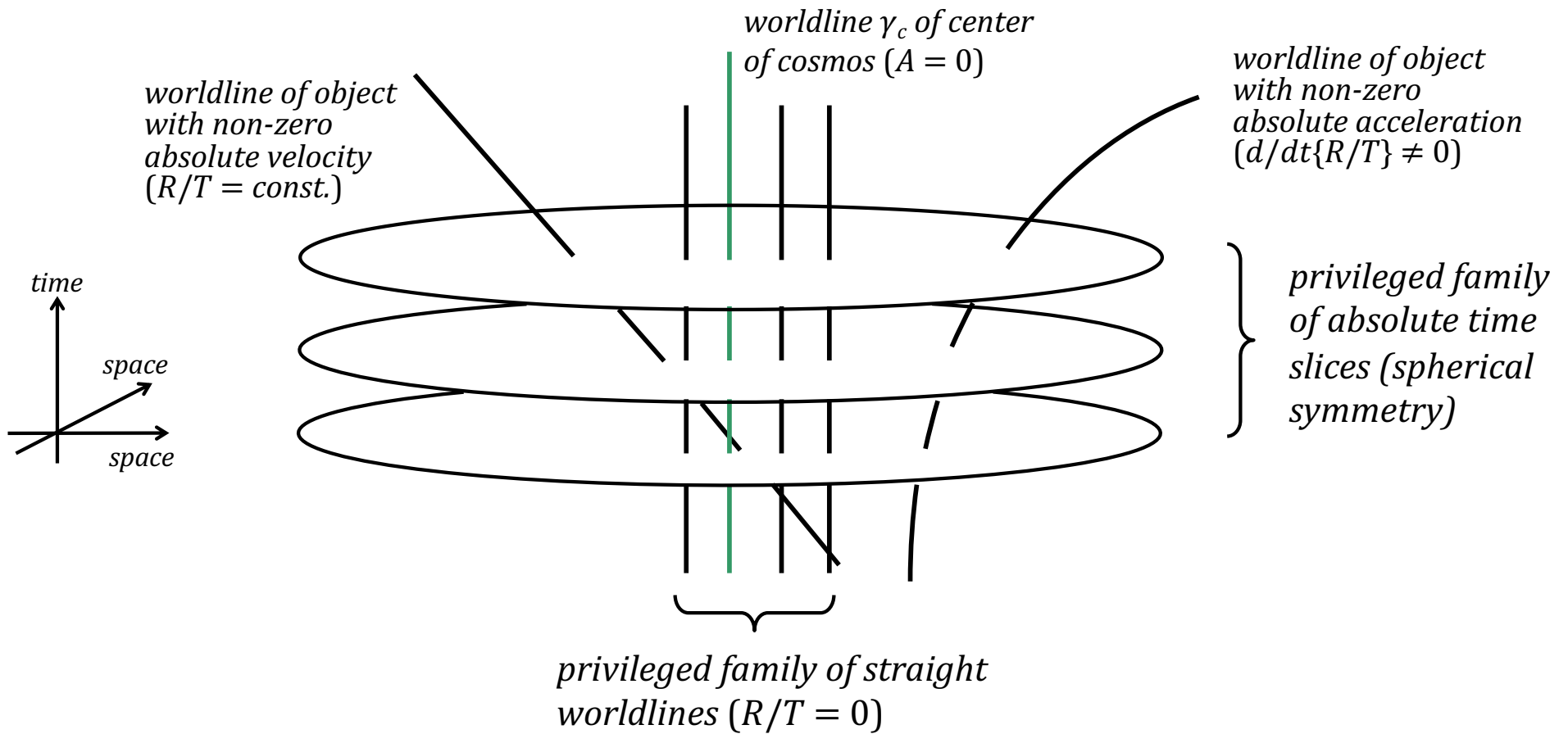
$$R(p, q) = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' - z)^2}$$

(A3) Any worldline  $\gamma$  through point  $p$  has a definite location  $A(\gamma, p)$ .

Consequence of (A3):

- Location is absolute.
- Define the worldline  $\gamma_c$  of the center of the cosmos by  $A(\gamma_c, p) = 0$ , for all points  $p$  on  $\gamma_c$ .

# Aristotelian Spacetime



1. Single, privileged inertial frame.
2. Position is absolute.
3. Velocity is absolute.
4. Acceleration is absolute.
5. Simultaneity is absolute.

- Newtonian, Galilean, and Aristotelian spacetimes support the Principle of Inertia (they all can distinguish straight worldlines from curved worldlines).
  - *Galilean spacetime does this minimally: no additional superfluous structure.*
  - *Newtonian and Aristotelian spacetimes have additional superfluous structure.*
- Are there classical spacetimes that do not have enough structure to distinguish straight worldlines from curved worldlines?



*James Clerk Maxwell*  
(1831-1879)

"Acceleration, like position and velocity, is a relative term and cannot be interpreted absolutely." (*Matter and Motion* 1877)

- *But*: With respect to Newton's Bucket Experiment...

"This concavity of the surface depends on the absolute motion of rotation of the water and not on its relative rotation."



- Absolute rotation but *no* absolute (linear) acceleration?

Maxwellian spacetime is a 4-dim collection of points such that:

(M1) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *temporal interval*  $T(p, q) = t' - t$ .

(M2) Between any two *simultaneous* points  $p_t, q_t$ , with coordinates  $(t, x, y, z)$  and  $(t, x', y', z')$ , there is a definite *Euclidean distance*,

$$R(p_t, q_t) = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' - z)^2}$$

(M3) Any worldline  $\gamma$  through point  $p$  has a definite twist  $\Omega(\gamma, p)$ .

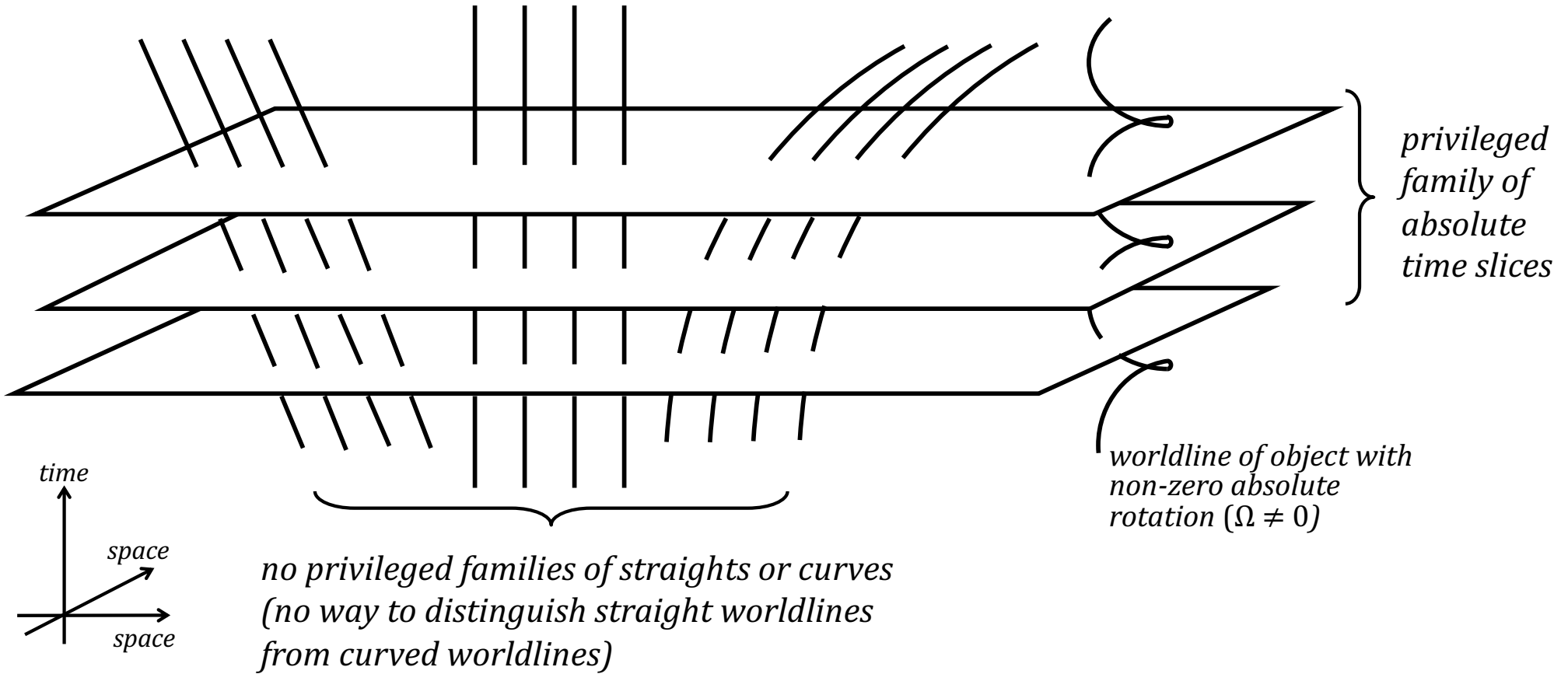
(M3) entails:

- *Linear* acceleration is no longer absolute!
  - *There is not enough structure in Maxwellian spacetime to distinguish straight worldlines from curved worldlines.*
- But rotation still is absolute!
  - *M3 allows us to tell when a worldline is "twisted".*

For worldline  $\gamma$  and point  $p$  on  $\gamma$ , the *absolute rotation* of  $\gamma$  with respect to  $p$  is given by  $\Omega(\gamma, p)$ .



# Maxwellian Spacetime



1. No inertial frames.
2. Velocity is relative.
3. Linear acceleration is relative.
4. Rotation is absolute.
5. Simultaneity is absolute.

Leibnizian spacetime is a 4-dim collection of points such that:

(L1) Between any two points  $p, q$ , with coordinates  $(t, x, y, z)$  and  $(t', x', y', z')$ , there is a definite *temporal interval*  $T(p, q) = t' - t$ .

(L2) Between any two *simultaneous* points  $p_t, q_t$ , with coordinates  $(t, x, y, z)$  and  $(t, x', y', z')$ , there is a definite *Euclidean distance*,

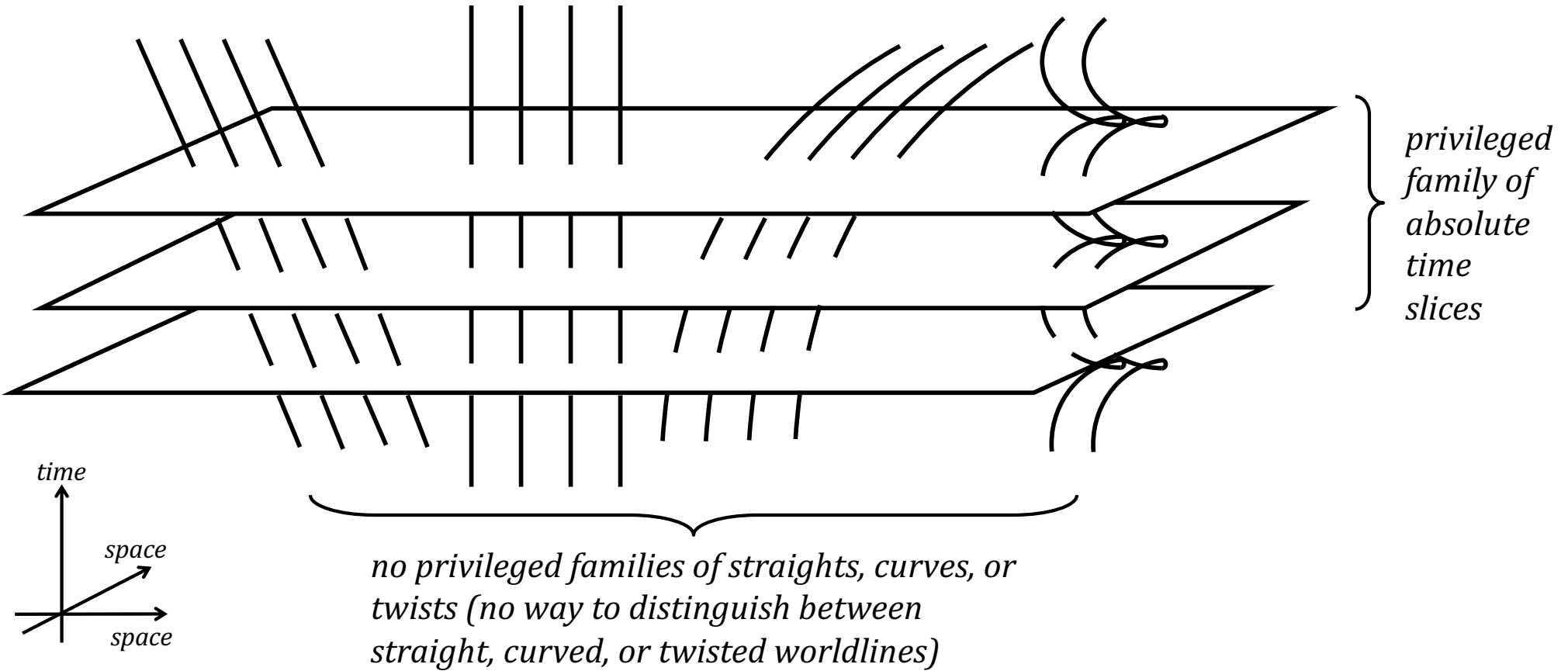
$$R(p_t, q_t) = \sqrt{(x' - x)^2 + (y' - y)^2 + (z' - z)^2}$$

- There's still an absolute temporal metric, and an absolute Euclidean spatial metric for the "instantaneous" 3-dim spaces.
  - *Space and time are absolute and Euclidean.*
- But: There's no rotation standard, and there's no acceleration standard.
- All motion is relative! (Relationism at last?)



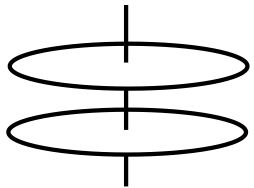
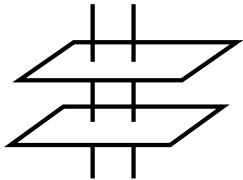
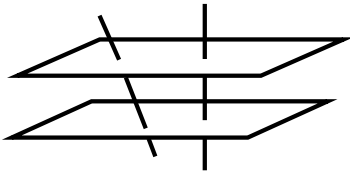
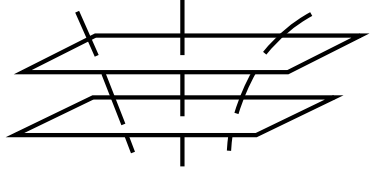
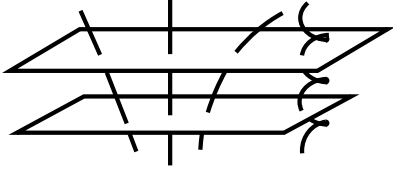
*Hmph! But I do concede that  
absolute acceleration exists...*

# Leibnizian Spacetime



1. No inertial frames.
2. Velocity is relative.
3. Acceleration is relative.
4. Rotation is relative.
5. Simultaneity is absolute.

# A Bestiary of Spacetimes

<i>Spacetime</i>	<i>Privileged Frames</i>	<i>Symmetries</i>	<i>Indistinguishable worldlines</i>
Aristotelian	Rigid Euclidean frame with position at origin, zero velocity, zero acceleration, zero rotation.	$x \rightarrow x' = \mathbf{R}x$ $t \rightarrow t' = t + \text{const.}$ <i>Rotate in space; translate in time.</i>	
Newtonian	Rigid Euclidean frame with zero velocity, zero acceleration, zero rotation.	$x \rightarrow x' = \mathbf{R}x + \text{const.}$ $t \rightarrow t' = t + \text{const.}$ <i>Rotate and translate in space; translate in time.</i>	
Galilean	Rigid Euclidean frames with zero acceleration, zero rotation.	$x \rightarrow x' = \mathbf{R}x + \mathbf{v}t + \text{const.}$ $t \rightarrow t' = t + \text{const.}$ <i>Rotate, translate and boost velocity in space; translate in time.</i>	
Maxwellian	Rigid Euclidean frames with zero rotation.	$x \rightarrow x' = \mathbf{R}x + \mathbf{a}(t)$ $t \rightarrow t' = t + \text{const.}$ <i>Rotate, translate, boost velocity and acceleration in space; translate in time.</i>	
Leibnizian	Rigid Euclidean frames.	$x \rightarrow x' = \mathbf{R}(t)x + \mathbf{a}(t)$ $t \rightarrow t' = t + \text{const.}$ <i>Rotate in space and time, translate and boost velocity and acceleration in space; translate in time.</i>	

Aristotelian

$$x \rightarrow x' = \mathbf{R}x$$

$$t \rightarrow t' = t + \text{const.}$$

Newtonian

$$x \rightarrow x' = \mathbf{R}x + \text{const.}$$

$$t \rightarrow t' = t + \text{const.}$$

Galilean

$$x \rightarrow x' = \mathbf{R}x + \mathbf{v}t + \text{const.}$$

$$t \rightarrow t' = t + \text{const.}$$

Maxwellian

$$x \rightarrow x' = \mathbf{R}x + \mathbf{a}(t)$$

$$t \rightarrow t' = t + \text{const.}$$

Leibnizian

$$x \rightarrow x' = \mathbf{R}(t)x + \mathbf{a}(t)$$

$$t \rightarrow t' = t + \text{const.}$$

- What transformations preserve Newton's 2nd Law:

$$F = md^2x/dt^2$$

- Answer: Aristotelian, Newtonian, and Galilean!

- Which means: The most general type of frames that cannot be experimentally distinguished by Newton's 2nd Law are Galilean frames!

- Consider what happens if we transform Newton's 2nd Law using Leibnizian transformations:

$$\begin{aligned}
 F &= m \frac{d^2}{dt^2} (\mathbf{R}(t)x + \mathbf{a}(t)) \\
 &= m(\ddot{\mathbf{R}}(t)x + 2\dot{\mathbf{R}}(t)\dot{x} + \mathbf{R}(t)\ddot{x} + \ddot{\mathbf{a}}(t))
 \end{aligned}$$

- The path of an object under no external forces is thus given by:

$$\ddot{x} + \mathbf{R}^{-1}(t)\ddot{\mathbf{R}}(t)x + \mathbf{R}^{-1}(t)\ddot{\mathbf{a}}(t) + 2\mathbf{R}^{-1}(t)\dot{\mathbf{R}}(t)\dot{x} = 0$$

"centrifugal inertial force"

"linear inertial force"

"coriolis inertial force"