

10. The Einstein Equations

Two requirements to geometrize the gravitational force in a relativistic spacetime:

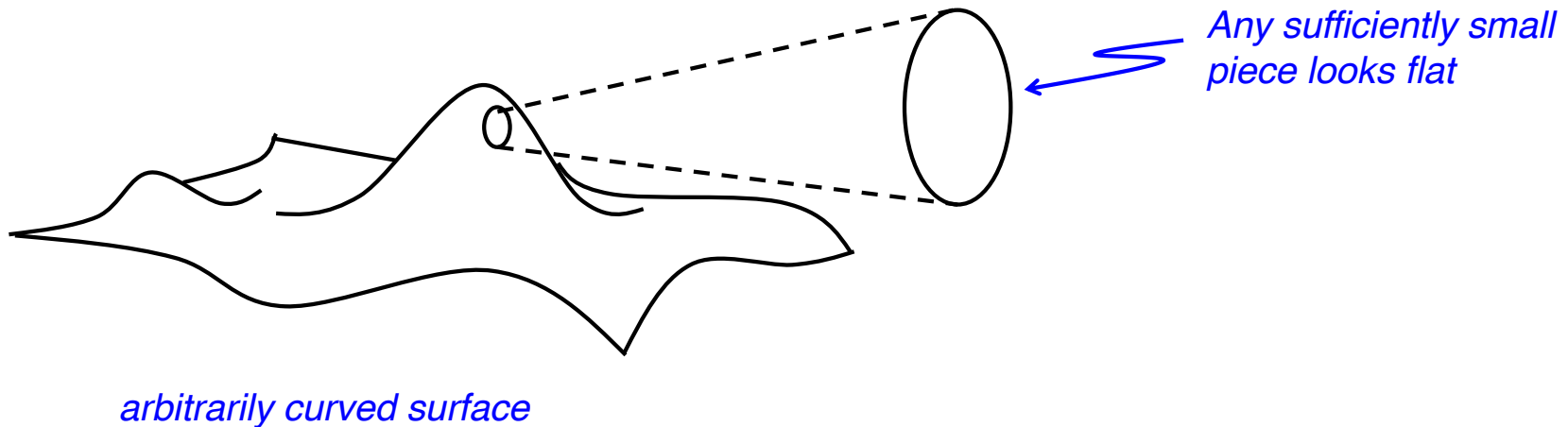
(I) Relativistic theory: Our theory must get everything right that special relativity gets right.

- Replace $ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu$ with $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$.


flat Minkowski metric


non-flat metric

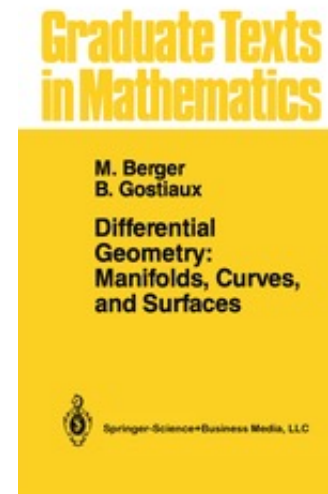
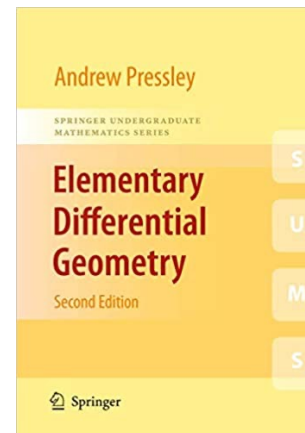
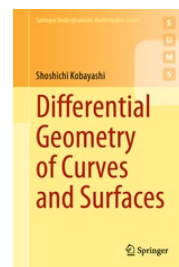
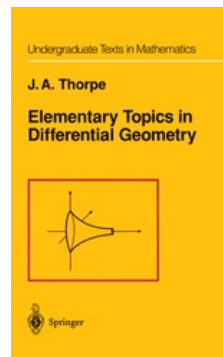
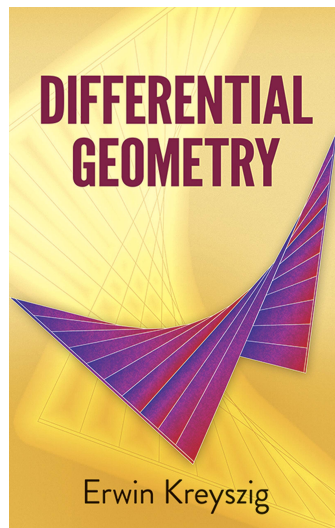
- Require $g_{\mu\nu}$ to reduce to $\eta_{\mu\nu}$ in small regions of spacetime.



(II) Geometrization: We want to construct a spacetime in which the equation for a straight line takes the form $(d^2x/dt^2 + \nabla\Phi) = 0$, in the *Newtonian limit*, where Φ is the Newtonian gravitational potential.

- *This means our theory will reproduce Newtonian gravity in the appropriate limit.*

To accomplish (II), Einstein had to learn differential geometry!



Two Facts from Differential Geometry

1. Geodesic Equation. The general equation for a straight line (a *geodesic*) in a curved space is given by,

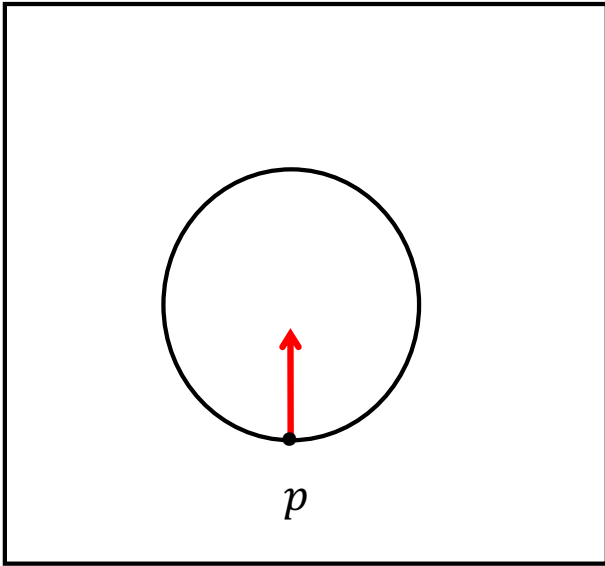
$$\frac{d^2 x^\mu}{dt^2} + \Gamma_{\nu\sigma}^\mu \frac{dx^\nu}{dt} \frac{dx^\sigma}{dt} = 0$$

- (a) The "correction factors" $\Gamma_{\nu\sigma}^\mu$ are (64!) functions that encode the curvature.
- (b) They depend explicitly on the metric $g_{\mu\nu}$.
- (c) $\Gamma_{\nu\sigma}^\mu = 0$ is a *sufficient*, but not *necessary*, condition for flatness.

2. Curvature Tensor. Gives us a *necessary* and *sufficient* condition for flatness.

Idea: Parallel-transport a tangent vector around a closed loop.

- *If the space is curved, it will come back pointing in a different direction!*

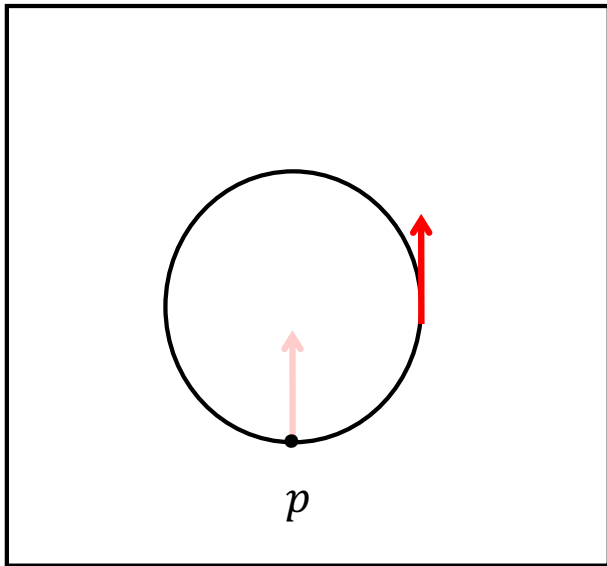


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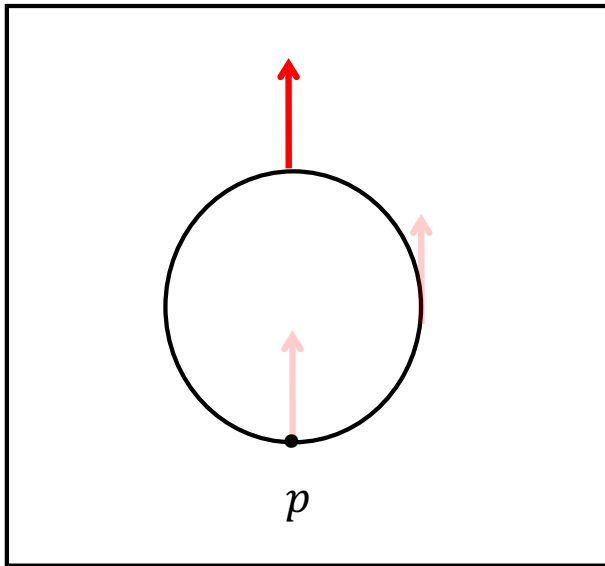


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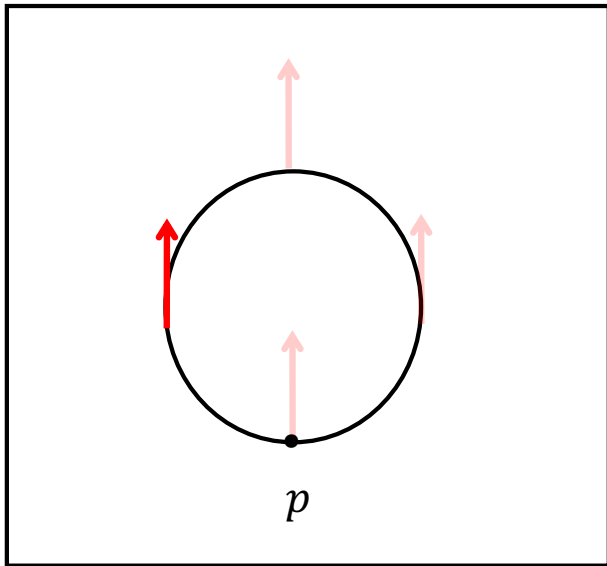


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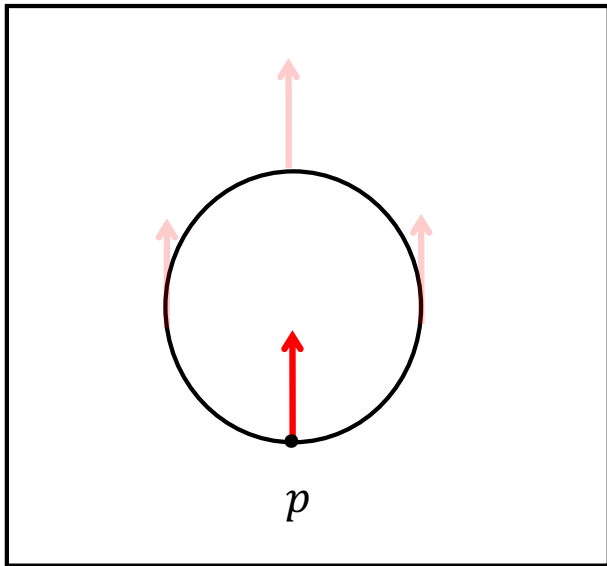


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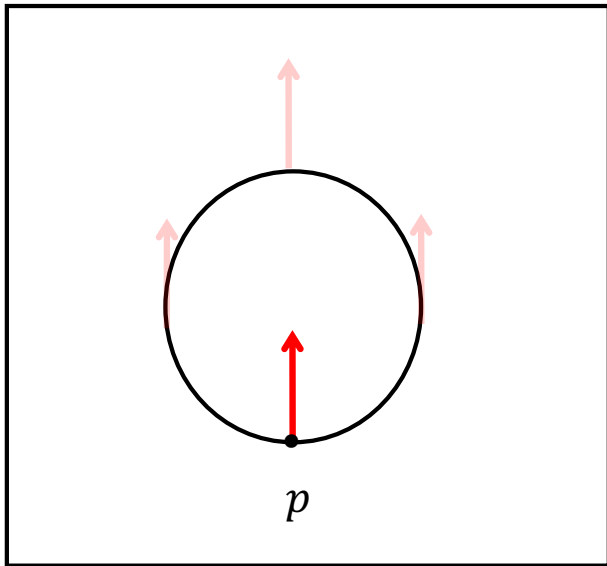


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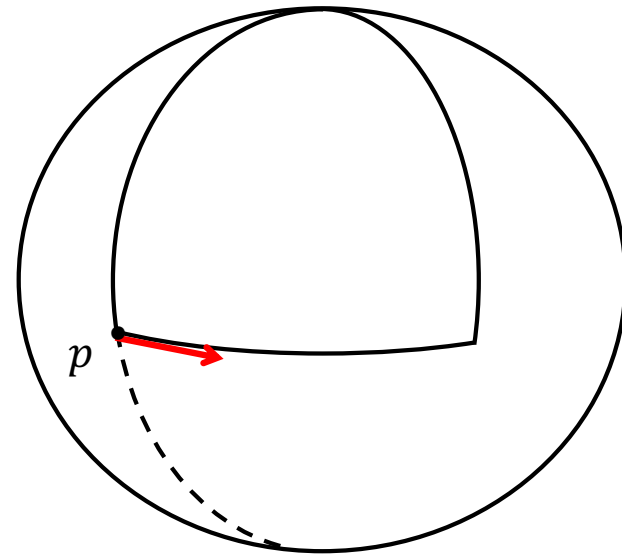
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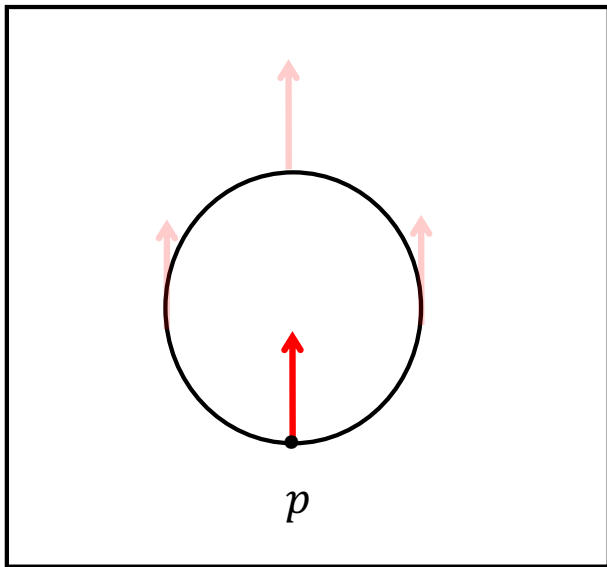


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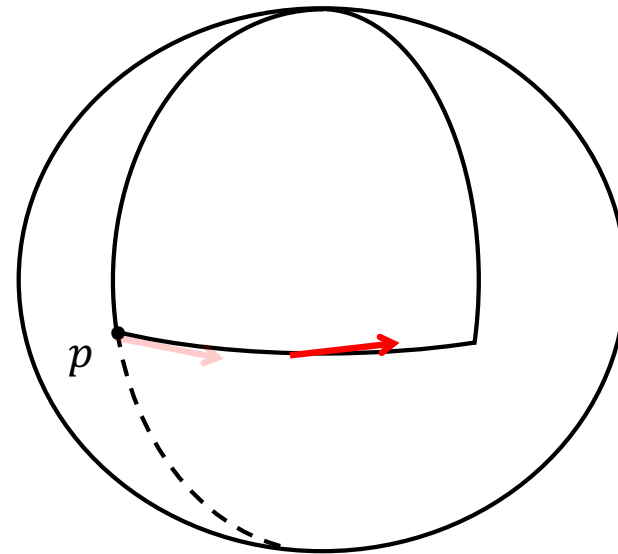
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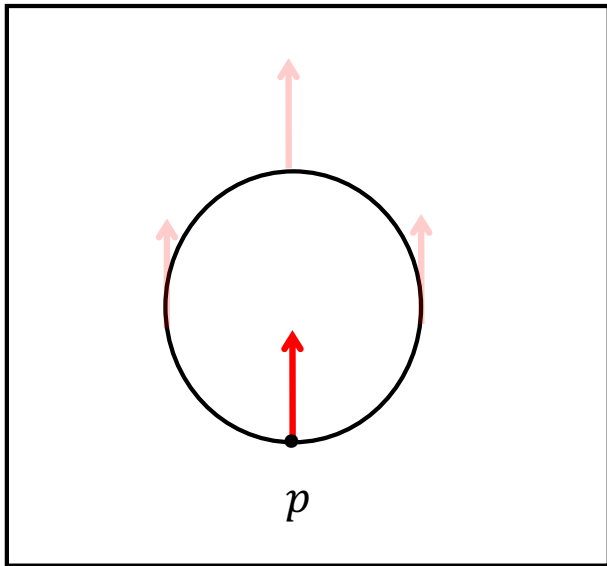


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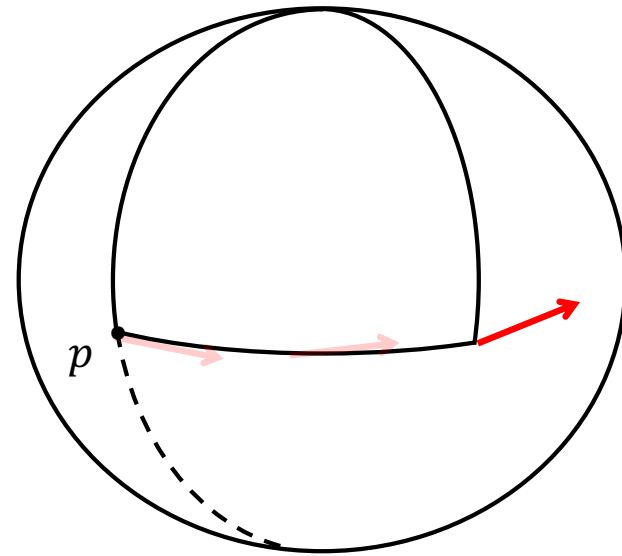
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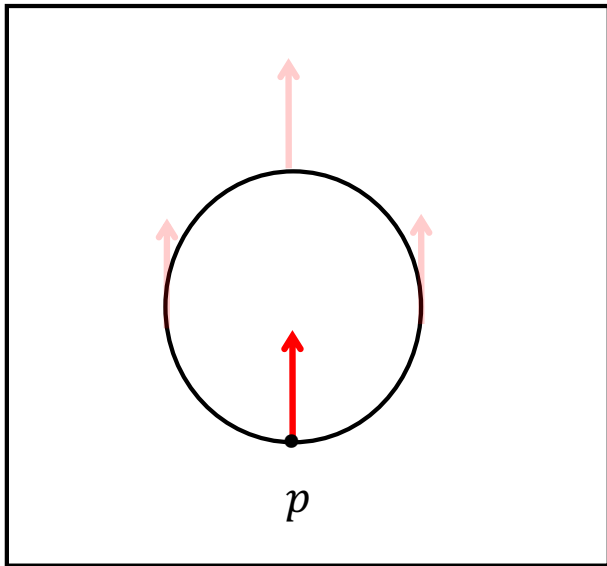


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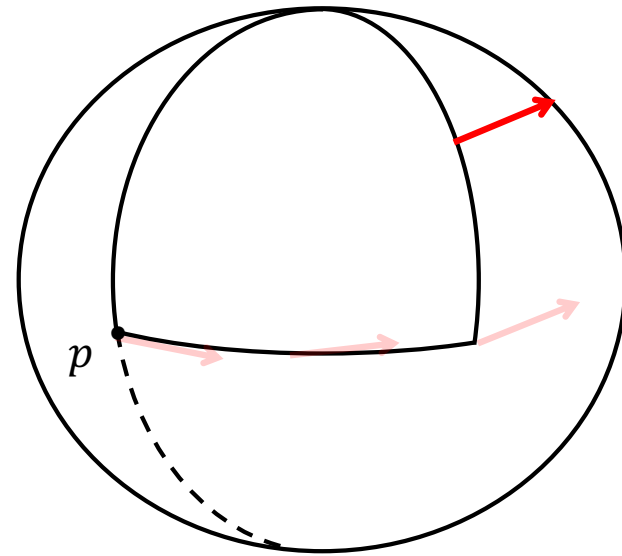
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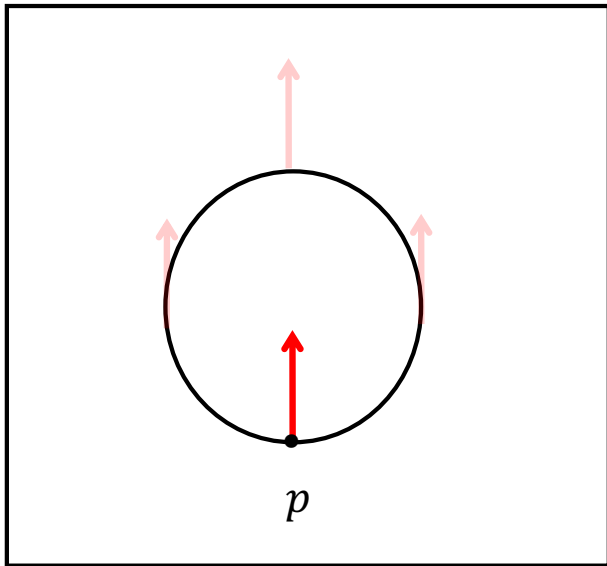


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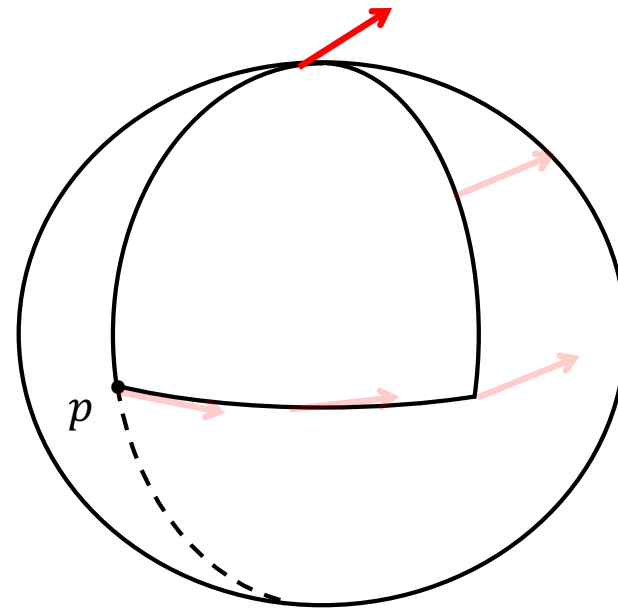
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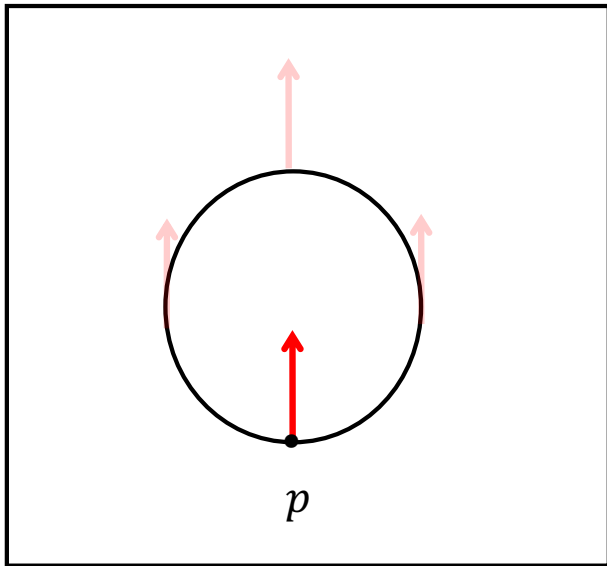


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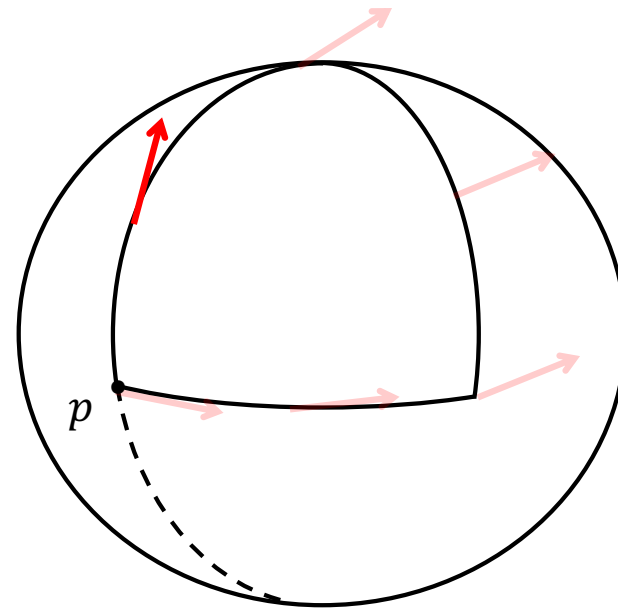
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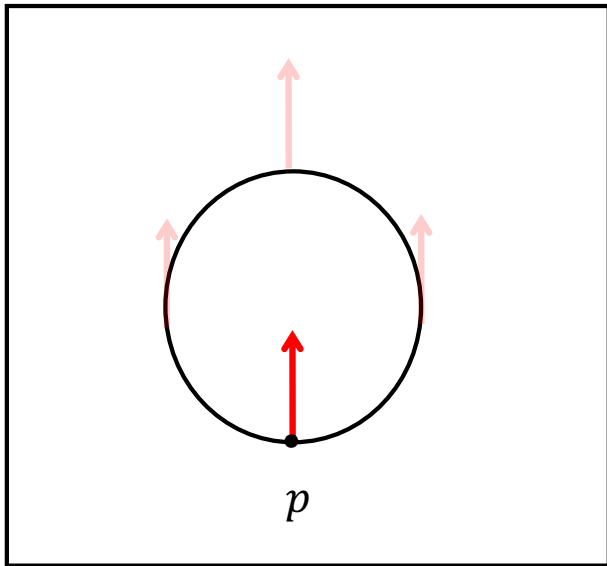


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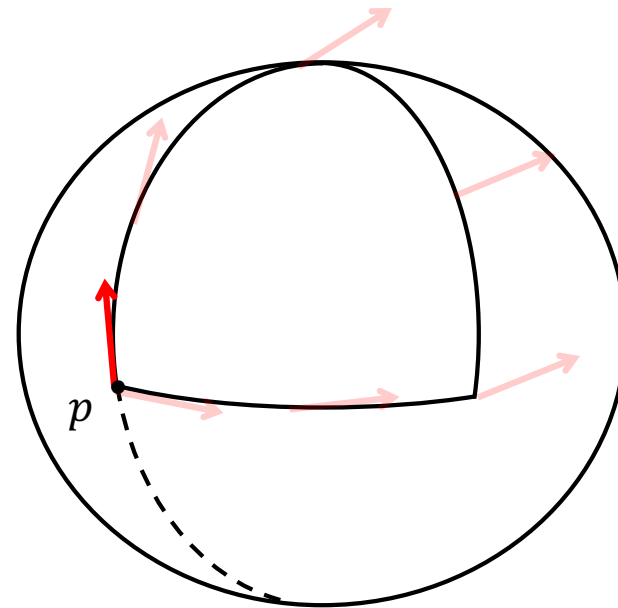
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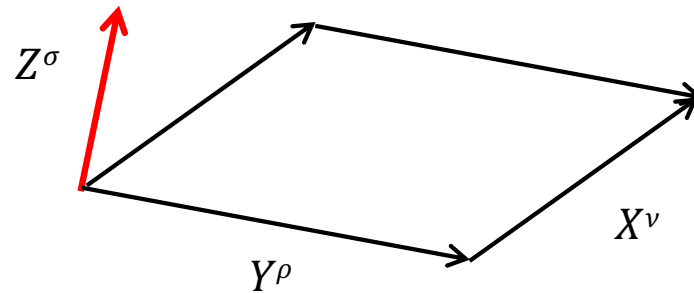


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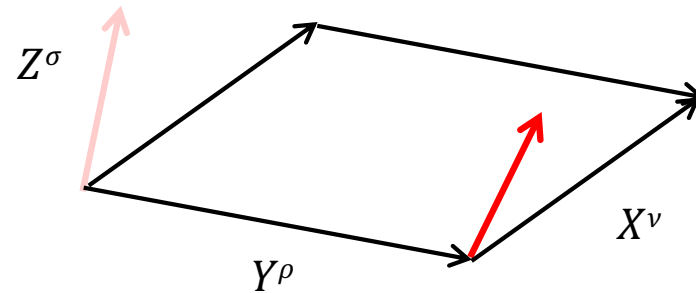


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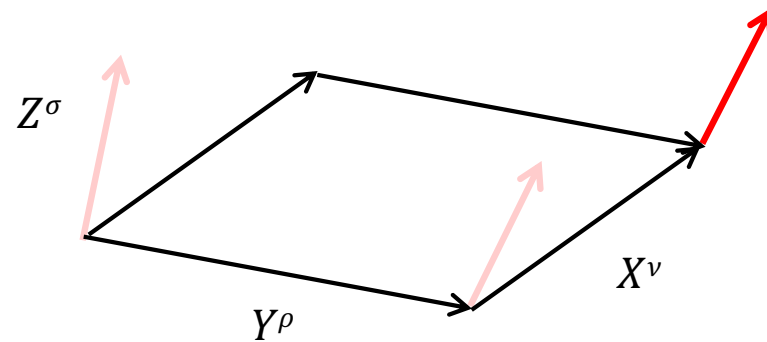
- Define a 4-indexed quantity, the *curvature tensor* $R^{\sigma}_{\mu\nu\rho}$, that measures this change.
- It acts on three vectors X^{ν} , Y^{ρ} , Z^{μ} and outputs the amount of change experienced by Z^{μ} upon parallel-transport around an infinitesimal curve defined by X^{ν} and Y^{ρ} :



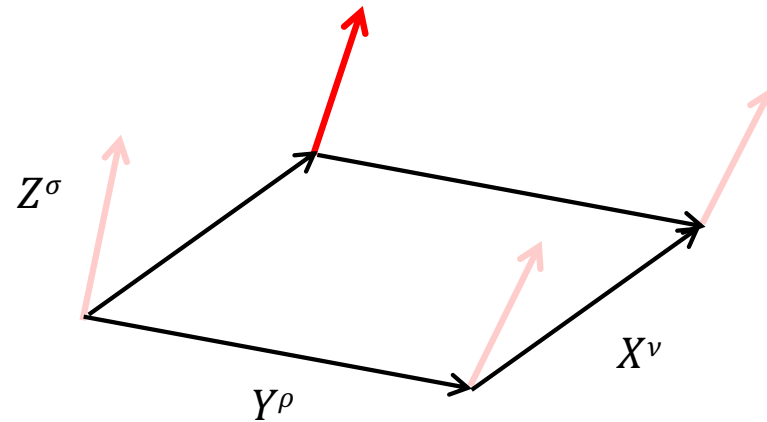
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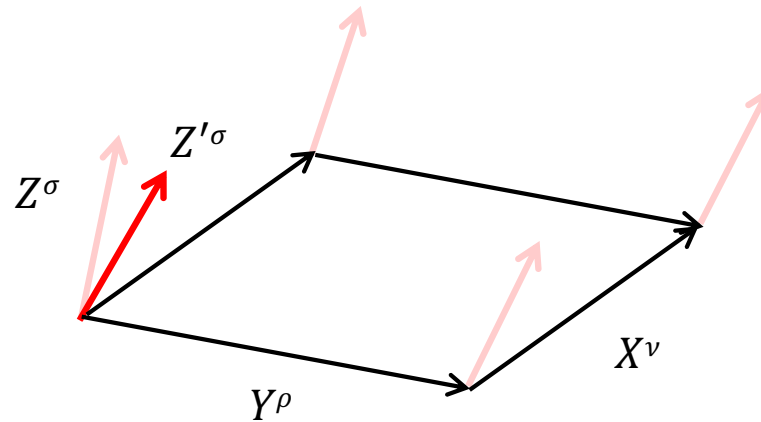
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$$\begin{aligned}
 R^\sigma_{\mu\nu\rho} X^\nu Y^\rho Z^\mu &= \delta Z^\sigma \\
 &= Z^\sigma - Z'^\sigma \\
 &= \text{change in } Z^\sigma \text{ upon parallel transport} \\
 &\quad \text{around loop defined by } X^\nu \text{ and } Y^\rho.
 \end{aligned}$$

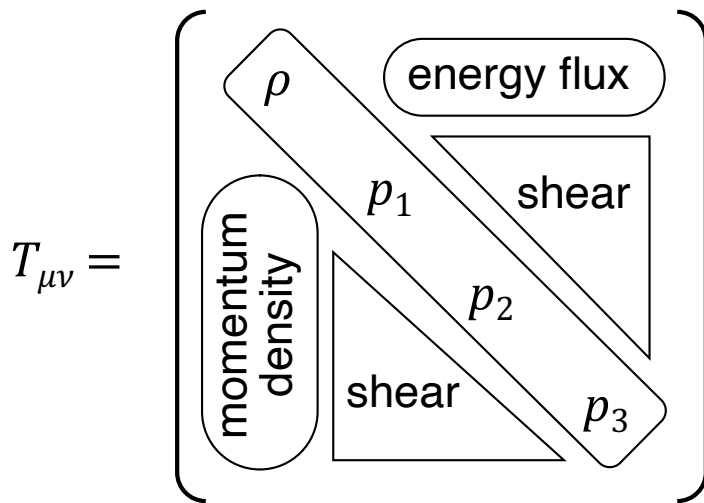
Properties of curvature tensor:

- $R^\sigma_{\mu\nu\rho} = 0$ if and only if the space is flat.
- $R^\sigma_{\mu\nu\rho}$ depends explicitly on the metric $g_{\mu\nu}$.

- To geometrize gravity, we need to relate the source of gravity (energy/mass) to the curvature tensor!
- We want there to be curvature in the presence of gravity, and no curvature in its absence.
- Require: (curvature of spacetime) \propto (matter density).

Can represent this mathematically by curvature tensor $R_{\mu\nu\rho}^{\sigma}$

Can represent this mathematically by "energy-momentum tensor" $T_{\mu\nu}$



$$T^{00} = \rho = \text{energy density}$$

$$T^{0i} = \text{energy flux components}$$

$$T^{i0} = \text{momentum density components}$$

$$T^{ii} = p_i = \text{pressure components}$$

$$T^{ij} = \text{shear components } (i \neq j)$$

$$i, j = 1, 2, 3$$

- Problem: The curvature tensor $R_{\mu\nu\rho}^{\sigma}$ has 4 indices and the energy-momentum tensor $T_{\mu\nu}$ has 2. Only tensors of the same "rank" can be equated!

- Eventual solution: Construct a 2-index tensor $G_{\mu\nu}$ out of $R^\sigma_{\mu\nu\rho}$ and set it equal to $T_{\mu\nu}$ with an appropriate proportionality constant.
- Result: The Einstein equations:

$$G_{\mu\nu}(g_{\mu\nu}) = \kappa T_{\mu\nu}$$

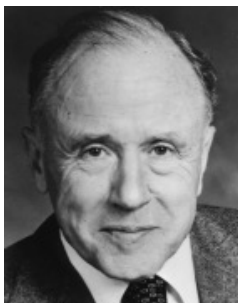
↗
"Einstein" tensor

Einstein tensor as a function of $g_{\mu\nu}$:

$$G_{\mu\nu}(g_{\mu\nu}) = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu}$$

↗ ↘
"Ricci" tensor "Ricci" scalar

- The constant $\kappa = 8\pi G$ ($G =$ Newtonian grav. constant) guarantees that these equations reproduce Newton's Law of Gravity in the Newtonian limit.
- Represent 16 differential equations of the metric $g_{\mu\nu}$, 6 of which are dependent on the rest; so 10 non-linear partial differential equations!
- What they mean:



John Wheeler
(1911-2008)

"...spacetime geometry tells matter how to move, matter tells spacetime how to curve"

- To solve the Einstein equations, one must make initial assumptions
 - *Either about spacetime geometry (e.g., isotropic; asymptotically flat, etc.).*
 - *Or about the matter distribution (e.g., evenly distributed, clumped in one spot, no negative energy, etc).*

A general relativistic spacetime = a 4-dim collection of points with the following additional structure: Between any two points, there is a spacetime interval given by $ds^2 = g_{\mu\nu}dx^\mu dx^\nu$, where $g_{\mu\nu}$ is a pseudo-Riemannian metric that satisfies the Einstein equations.

"reduces to the Minkowski metric at any point"