Emergent Spacetime and Structural Realism

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Prospects for modeling spacetime as a phenomenon that emerges in the low-energy limit of a quantum liquid.

- 1. EFTs in Condensed Matter Systems
- 2. Superfluid Helium
- **3. Emergence and Emergent Spacetime**
- 4. Universality, Topology, and Dynamical Structure

1. EFTs in Condensed Matter Systems

Highly-correlated condensed matter system

<u>Non-relativistic</u> many-body quantum system that displays macroscopic quantum effects

- superfluids
- superconductors
- Bose-Einstein condensates
- quantum Hall liquids

Effective Field Theory of condensed matter system

Theory of low-energy dynamics of system: describes states with energy close to zero

- bosonic collective modes of ground state
- fermionic excitations above ground state ("quasiparticles")
- topological defects ("vortices")

1. EFTs in Condensed Matter Systems

How to Construct a Condensed Matter EFT Take Low-energy "limit":

• Expand Lagrangian in small fluctuations in field variables about ground state and integrate out high-energy fluctuations. $({}^{4}He)$

OR

• Linearize the energy about the values where it vanishes and then construct the corresponding Hamiltonian. $({}^{3}He-A)$

For fermionic liquids, the type of EFT that results is ultimately based on *topological* properties of the ground state of system, as opposed to its *symmetries*.

2. "Acoustic" Spacetimes and Superfluid ⁴He

- liquid consisting of many ${}^{4}He$ atoms, all phases alligned
- model ground state as single quantum particle: $\varphi_0 = \rho_0 e^{i\theta}$

$$\mathcal{L}_{4He} = i arphi^{\dagger} \partial_t arphi - rac{1}{2m} \partial_i arphi^{\dagger} \partial_i arphi + \mu arphi^{\dagger} arphi - lpha^2 (arphi^{\dagger} arphi)^2 \lambda^{\dagger} + rac{1}{2m} \partial_i arphi^{\dagger} \partial_i arphi + rac{1}{2m} \partial_i arphi^{\dagger} \partial_i arphi^{\dagger} \partial_i arphi^{\dagger} \partial_i arphi + rac{1}{2m} \partial_i arphi^{\dagger} \partial_i arph^{\dagger} \partial_i arphi^{\dagger} \partial_i arphi^{\dagger} \partial_i arphi^{\dagger} \partial_i$$

kinetic energy

conservation of particle#

SSB potential Non-relativistic Lagrangian for Superfluid ⁴He

Low-energy limit:

- Let $\varphi = \rho e^{i\theta}, \rho = \overline{\rho_0 + \delta\rho}, \theta = \theta_0 + \delta\theta$
- Integrate out high-energy fluctuations $\delta \rho$

$$\mathcal{L}_{4He} = \mathcal{L}_0[
ho_0,\, heta_0] + \mathcal{L}'_{4He}[\delta heta]$$

ground state

low-energy fluctuations above ground state (EFT) 2. "Acoustic" Spacetimes and Superfluid ⁴He

$$\mathcal{L'}_{4He} = rac{1}{4lpha^2} (\partial_t heta + v_i \partial_i heta)^2 - rac{
ho_0}{2m} (\partial_i heta)^2 \qquad i = 1, 2, 3$$

 $EFT \; for \; Superfluid \; {}^4He \ egin{pmatrix} to \; 2nd \; order \; in \; \delta heta : \ v_i = (1/m) \partial_i heta \end{pmatrix}$



 $\mathcal{L}'_{4He} = rac{1}{2} \sqrt{g} \; g^{\mu
u} \partial_{\mu} \theta \partial_{
u} \theta \qquad \mu, \;
u = 0, \, 1, \, 2, \, 4$

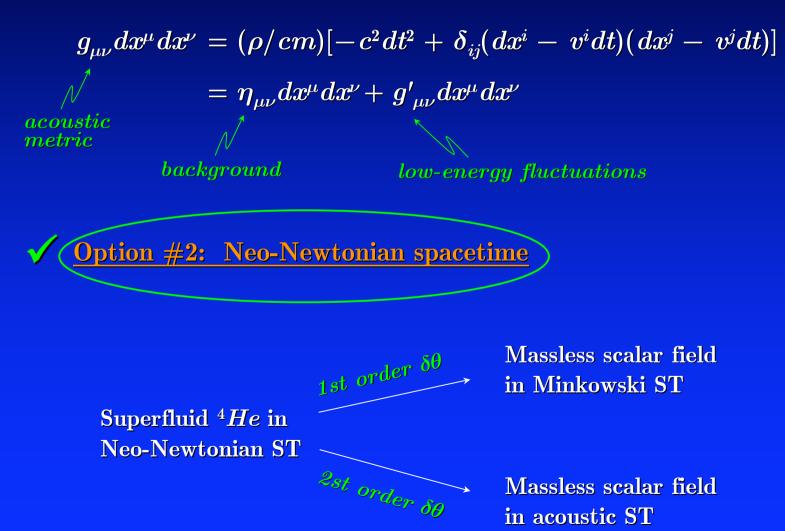
Massless scalar field in curved spacetime!

 $egin{aligned} g_{\mu
u}dx^{\mu}dx^{
u} &= (
ho/cm)[-c^2dt^2+\delta_{ij}(dx^i-v^idt)(dx^j-v^jdt)]\ c^2 &\equiv 2lpha^2
ho/m \end{aligned}$

Can now model black hole physics:

speed of light = speed of low-energy ocillations (i.e., "sound" modes)
<u>Hence</u>: "acoustic" spacetimes and "acoustic" black holes

- 2. "Acoustic" Spacetimes and Superfluid ⁴He
- (i) What is the kinematic background of acoustic spacetimes?
 Option #1: Minkowski spacetime



- 2. "Acoustic" Spacetimes and Superfluid ${}^{4}He$
- To what extent are "acoustic" spacetimes analogues of GR (ii) spacetimes?

Einstein equations cannot
be derived from ${}^{4}He$ EFT. \Rightarrow $\begin{bmatrix} Not \ dynamical \ analogues! \end{bmatrix}$



Kinematic analogues of GR?

... the features of general relativity that one typically captures in an "analogue model" are the *kinematic* features that have to do with how fields (classical or quantum) are defined on curved spacetime, and the sine qua non of any analogue model is the existence of some "effective metric" that captures the notion of the curved spacetimes that arise in general relativity. (Barceló, Liberati, Visser 2005, pg. 10.)

- 2. "Acoustic" Spacetimes and Superfluid ${}^{4}He$
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 \Rightarrow [Not dynamical analogues!]

Kinematic analogues of GR?

The acoustic analogue for black-hole physics accurately reflects half of general relativity -- the kinematics due to the fact that general relativity takes place in a Lorentzian spacetime. The aspect of general relativity that does not carry over to the acoustic model is the dynamics -- the Einstein equations. Thus the acoustic model provides a very concrete and specific model for separating the kinematic aspects of general relativity from the dynamic *aspects.* (Visser 1998, pg. 1790.)

- 2. "Acoustic" Spacetimes and Superfluid ${}^{4}He$
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Einstein equations cannot be derived from ⁴*He* EFT.

 \Rightarrow [Not dynamical analogues!]

Kinematic analogues of GR?

- If kinematics of GR = Minkowski ST, then No!
- Kinematics of GR = ?

Some features that one normally thinks of as intrinsically aspects of gravity, both at the classical and semiclassical levels (such as horizons and Hawking radiation), can in the context of acoustic manifolds be instead seen to be rather generic features of curved spacetimes and quantum field theory in curved spacetimes, that have nothing to do with gravity per se. (Barceló, Liberati, Sonego, Visser 2004, pg. 2.)

2. Standard Model and Superfluid ³He-A

- Liquid of many ³He Cooper Pairs, all phases alligned
- Degrees of freedom: $S_z = 0, \pm 1; l_z = 0, \pm 1$
- A-phase: no $S_z = 0$ substates, $\hat{d} \mid\mid \hat{l}$,

$$H_{\!_{3He\text{-}A}} = \chi^{\dagger}_{lphaeta} \{(arepsilon-\mu)\sigma^{}_3 + \,ec{V}_{lphaeta}(\,\widehat{l}\!,\,\widehat{d}\!,\,ec{k}\!)\,\}\chi_{lphaeta}$$

$$E(ec{k})=0,~~ ext{for}~2~ ext{values}~ ext{of}~ec{k}$$

<u>Low-energy limit</u>

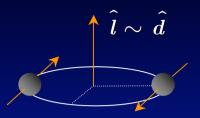
• Expand $E(\vec{k})$ to 2nd order about zero points:

$$E^2(\,\overline{k})pprox\, g_{ij}(k_i^{}-\,qA_i^{})(k_j^{}-\,qA_j^{}) \qquad \qquad g_{ij}^{}\sim l_i^{}l_j^{}\,, A_i^{}\sim l_i^{}$$

• Construct $\mathcal{L}'_{_{3He-A}}$

$$egin{aligned} \mathcal{L}'_{{}_{3He}\!-\!A} &= \overline{\Psi} g_{\mu
u} \gamma^
u (\partial_\mu - q A_\mu) \Psi \ g_{\mu
u} &\sim (l_i,\,v_i),\, A_0 \sim l_i v_i \end{aligned}$$

Potential field A_{μ} interacting with matter field Ψ in curved spacetime



(Volovik 2003)

- 2. Standard Model and Superfluid ${}^{3}He-A$
- "Induced QED" (Zeldovich 1967): expand to 2nd order in fluctuations in A_{μ}

$${\cal L}^{\,\prime}_{_{3He-A}} = \overline{\Psi} g_{\mu
u} \gamma^
u (\partial_\mu - q A_\mu) \Psi + rac{1}{4\kappa^2} \sqrt{-g} \, g^{\mu
u} F_{\mulpha} F_{
ueta} \, ,$$

(3+1)-dim QED in curved spacetime

- Extendable in-principle to SU(n) gauge fields \Rightarrow Standard Model
- Similar treatment of effective metric ("induced gravity" Sakharov 1967) fails to reproduce Einstein-Hilbert term

<u>Interpretation</u>

ground state of ${}^{3}He-A \iff \text{spacetime } (g_{\mu\nu})$ low-energy fluctuations (quasiparticles, collective modes) $\iff \text{matter fields, potential fields } (\Psi, A_{\mu})$ "induced" vacuum corrections to

interactions between Ψ and A_{μ}



gauge fields $(F_{\mu\nu})$

3. Emergence...

<u>Claim:</u> Novel phenomena (fields, particles, symmetries, spacetime, etc.) emerge in low-energy limit of certain condensed matter systems

Emergence in the low-energy limit

(1) Distinct from emergence via symmetry breaking.



(2) Epistemological Emergence

- Unpredictability
- Irreducibility (and/or unexplainability)
- Causal efficacy

3. ... and Emergent Spacetime

Spacetime as emergent...

(a) ... in superfluid ⁴He

- *Motivation*: Spacetime as (some aspect of) solutions to Einstein equations in GR.
- Prospects: Limited. ⁴He EFT lacks GR dynamics/kinematics.

(b) ... in superfluid ³He-A

- *Motivation*: Spacetime as ground state for QFTs of matter, gauge, and metric fields.
- *Prospects*: Still limited. ³*He-A* EFT does not fully recover GR.

<u>Research programme</u>:

Determine appropriate condensed matter system that produces relevant matter, gauge and metric fields in low-energy limit. (Volovik 2003, Wen 2004.) 3. Universality, Topology, ...

Why does ³He-A reproduce the Standard Model?

- Universal property = property stable under perturbations
- Universality class = class of systems/states characterized by common universal properties
- Universality class of the ground state of a system determines it's lowenergy dynamics (*i.e.*, the corresponding EFT)

• Universal properties of ground state are "generic" properties of EFT:

- how system behaves at a phase transition
- whether theory is renormalizable
- symmetries of low-energy fluctuations (i.e., symmetries of EFT)

3. Universality, Topology, ...

Why does ³He-A reproduce the Standard Model?

- Universality classes of fermionic ground states are characterized by momentum space topology.
 - stable regions in k-space where quasiparticle energies $\rightarrow 0$
- ³*He-A* and the Standard Model have ground states characterized by the *same* momentum space topology.

- stable point defects ("Fermi points")

<u>Thus</u>: ³He-A and the Standard Model possess the same universal properties.

3. ... and Dynamical Structure



Irrespective of microscopic details:

- Standard Model
- ³*He*-*A*
- any condensed matter system with "Fermi points"

Same dynamical structure

<u>Structural Realist interpretation of spacetime:</u>

• spacetime as a universal property = generic property of an EFT universality class

<u>Qualification</u>: Universality class that best describes spacetime structure still unknown